

A status update on the
neutrino mechanism of core-collapse supernovae

J. Nordhaus
Center for Computational Relativity and Gravitation
Rochester Institute of Technology (USA)

Introduction:

What questions do we aim to answer?
The core-collapse puzzle and its history.

Physical Processes:

Core collapse.
The stalled shock and its revival.

The Delayed-Neutrino Mechanism:

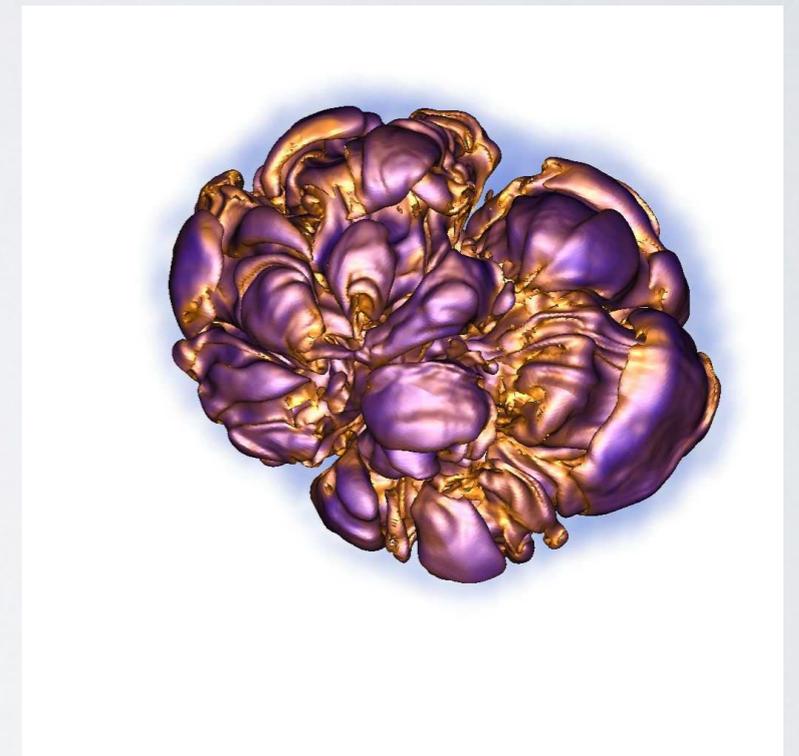
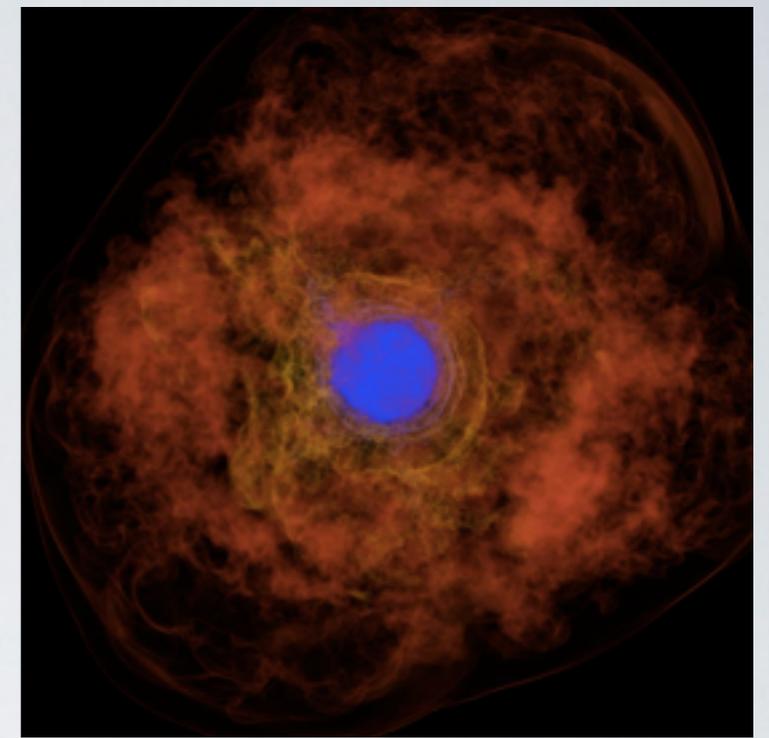
The status of current simulations.

Alternative Ideas:

Jets / nuclear burning / GR / MHD?

Pulsar Kicks:

Recoil from core-collapse explosions.



Collaborators: T. Brandt (IAS), M. Rantsiou (MPA), C. Ott (Caltech), E. Livne (Hebrew Univ.), A. Almgren (LBL), J. Bell (LBL), A. Burrows (Princeton), O. Papish (Technion), N. Soker (Technion)

“Every passing hour brings the Solar System 43,000 miles closer to Globular Cluster M13 in Hercules - and still there are some misfits who insist there is no such thing as progress” - Ransom K. Fern

Kurt Vonnegut
The Sirens Of Titan

Core-collapse supernovae background

Progenitors: Massive stars $8 - 50M_{\odot}$

Energies:

Neutron star binding energy: $\sim 3 \times 10^{53}$ erg

Kinetic energy of ejecta: $\sim 10^{51}$ erg

Light-curve energy: $\sim 10^{49}$ erg

Timescales:

Core collapse: ~ 500 ms

Post-bounce time to explosion: ≤ 1 s

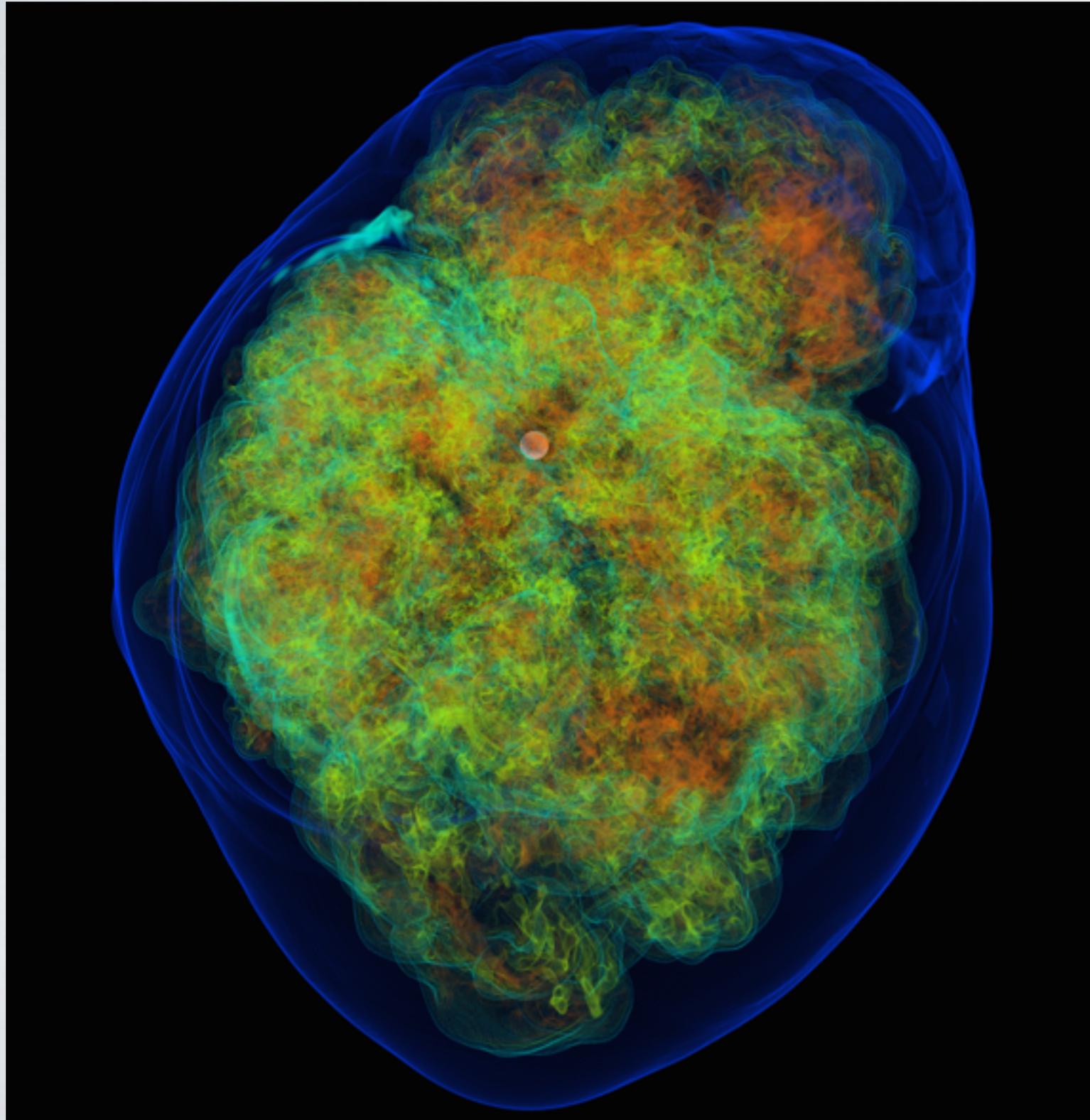
Core neutrino cooling time: ~ 10 s

Products: Blast wave, heavy-elements, compact object remnants.

What questions do we aim to answer?

1. Canonical explosion energy of 1 Bethe.
2. Neutron star mass distribution.
3. Nucleosynthetic yields and distributions.
4. Pulsar Kicks.

What physics is necessary?



Potentially Important Ingredients

- ▶ Gravity
- ▶ Neutrino Heating
- ▶ Turbulence/Convection
and Shock Instabilities
- ▶ Rotation
- ▶ Magnetic fields
- ▶ Nucleosynthesis
- ▶ General Relativity

Multi-dimensional effects
important!

Goal: 3D models with
sufficient realism that produce
canonical SN explosions

The core-collapse puzzle and its history

Chandrasekhar (1935): Iron core must collapse if $M_c \geq 1.4M_\odot$

The core-collapse puzzle and its history

Chandrasekhar (1935): Iron core must collapse if $M_c \geq 1.4M_\odot$

Burbidge, Burbidge, Fowler, Hoyle (1957):

Photo-disintegration of iron removes:

$$1.65 \times 10^{18} \text{ erg g}^{-1} \gg U_{\text{therm}} \sim 3 \times 10^{17} \text{ erg g}^{-1}$$

Nuclear burning on infall.

The core-collapse puzzle and its history

Chandrasekhar (1935): Iron core must collapse if $M_c \geq 1.4M_\odot$

Burbidge, Burbidge, Fowler, Hoyle (1957):

Photo-disintegration of iron removes:

$$1.65 \times 10^{18} \text{ erg g}^{-1} \gg U_{\text{therm}} \sim 3 \times 10^{17} \text{ erg g}^{-1}$$

Nuclear burning on infall.

Colgate & White (1961-1966):

Electron capture, nuclear dissociation initiate dynamical collapse.

Collapse halted at nuclear densities, shock wave forms as core matter is suddenly decelerated.

Nuclear burning won't work.

Neutrinos proposed as an energy transport mechanism.

The core-collapse puzzle and its history

Chandrasekhar (1935): Iron core must collapse if $M_c \geq 1.4M_\odot$

Burbidge, Burbidge, Fowler, Hoyle (1957):

Photo-disintegration of iron removes:

$$1.65 \times 10^{18} \text{ erg g}^{-1} \gg U_{\text{therm}} \sim 3 \times 10^{17} \text{ erg g}^{-1}$$

Nuclear burning on infall.

Colgate & White (1961-1966):

Electron capture, nuclear dissociation initiate dynamical collapse.

Collapse halted at nuclear densities, shock wave forms as core matter is suddenly decelerated.

Nuclear burning won't work.

Neutrinos proposed as an energy transport mechanism.

Early simulations: Bounce shock stalls at $\sim 100 - 200$ km.

Sapped of pressure by electron capture, neutrino losses.

The core-collapse puzzle and its history

~ **1970 - present:** Focus on the neutrinos.

The core-collapse puzzle and its history

~ **1970 - present:** Focus on the neutrinos.

Arnett (1967, 1977), Sato (1975):

Lower densities implies neutrinos are less trapped: $\rho \sim 10^{13} \text{ g cm}^{-3}$
Burst of flux over $\sim 100 \text{ ms}$, radiation pressure ejects envelope?

The core-collapse puzzle and its history

~ 1970 - present: Focus on the neutrinos.

Arnett (1967, 1977), Sato (1975):

Lower densities implies neutrinos are less trapped: $\rho \sim 10^{13} \text{ g cm}^{-3}$
Burst of flux over ~ 100 ms, radiation pressure ejects envelope?

Bethe & Wilson (1985):

Nuclei do not dissociate on infall. Bounce at $\rho \sim 2.6 \times 10^{14} \text{ g cm}^{-3}$
Shock energized with pdV work, not neutrino pressure.

The core-collapse puzzle and its history

~ 1970 - present: Focus on the neutrinos.

Arnett (1967, 1977), Sato (1975):

Lower densities implies neutrinos are less trapped: $\rho \sim 10^{13} \text{ g cm}^{-3}$
Burst of flux over $\sim 100 \text{ ms}$, radiation pressure ejects envelope?

Bethe & Wilson (1985):

Nuclei do not dissociate on infall. Bounce at $\rho \sim 2.6 \times 10^{14} \text{ g cm}^{-3}$
Shock energized with pdV work, not neutrino pressure.

1D and 2D simulations (1995 - 2010):

Spherically symmetric: do not explode.

Axisymmetric: some explosions for weakly-bound atmospheres.

The core-collapse puzzle and its history

~ 1970 - present: Focus on the neutrinos.

Arnett (1967, 1977), Sato (1975):

Lower densities implies neutrinos are less trapped: $\rho \sim 10^{13} \text{ g cm}^{-3}$
Burst of flux over ~ 100 ms, radiation pressure ejects envelope?

Bethe & Wilson (1985):

Nuclei do not dissociate on infall. Bounce at $\rho \sim 2.6 \times 10^{14} \text{ g cm}^{-3}$
Shock energized with pdV work, not neutrino pressure.

1D and 2D simulations (1995 - 2010):

Spherically symmetric: do not explode.

Axisymmetric: some explosions for weakly-bound atmospheres.

3D simulations (2010 -):

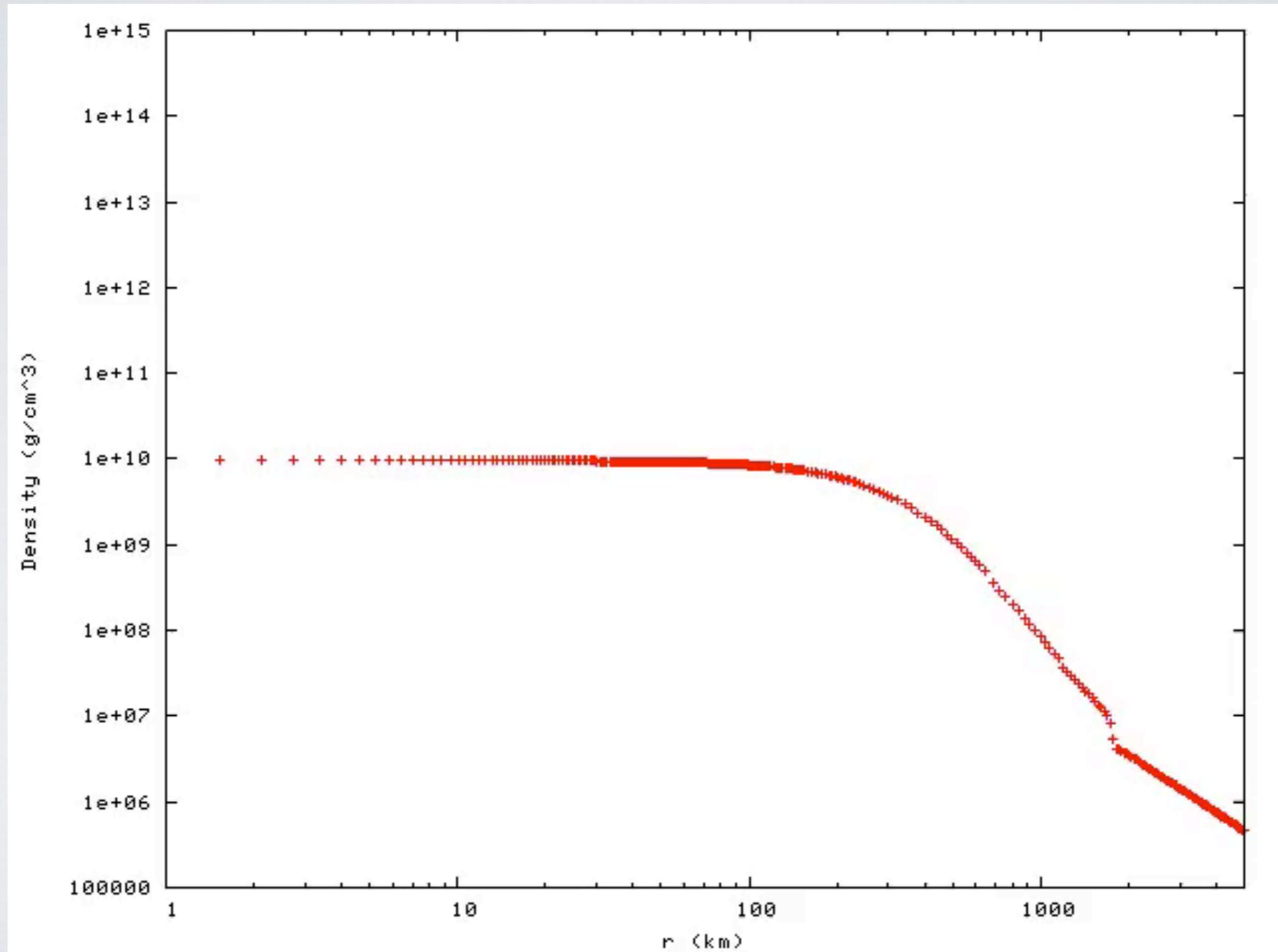
Detailed neutrino transport

Some weak explosions.

MHD, GR, other secondary physics.

No realistic simulation has produced a supernova with canonical energy.

The core-collapse puzzle and its history



Physics of core collapse

Massive iron core electrons become relativistically degenerate.

Adiabatic exponent of $\gamma \leq \frac{4}{3}$ implies unstable to dynamical collapse.

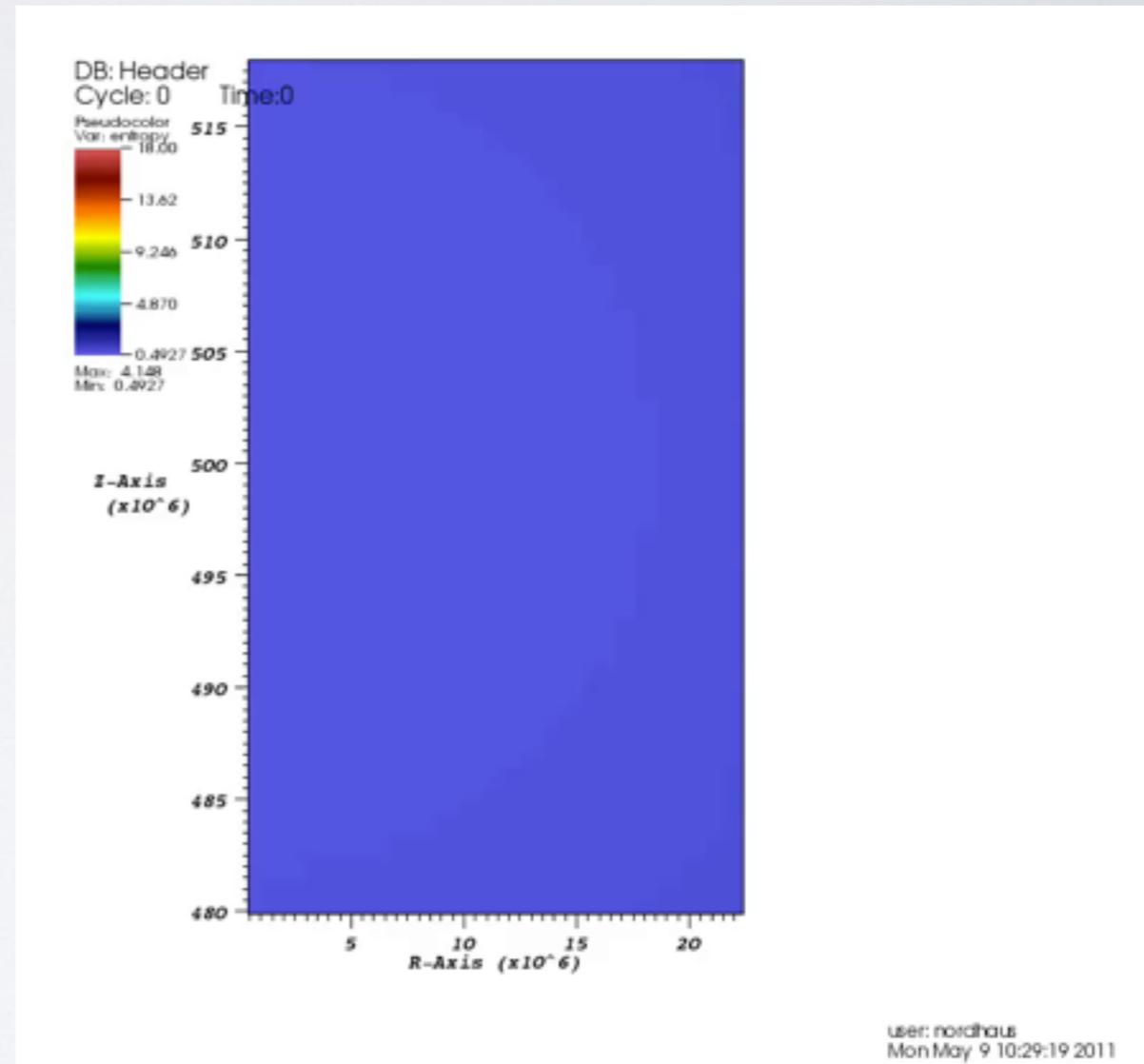
Collapse halted when equation of state stiffens due to nucleons.

Inner core:

Remains in sonic contact, rebounds as a unit.

Bounce dissociates nuclei.

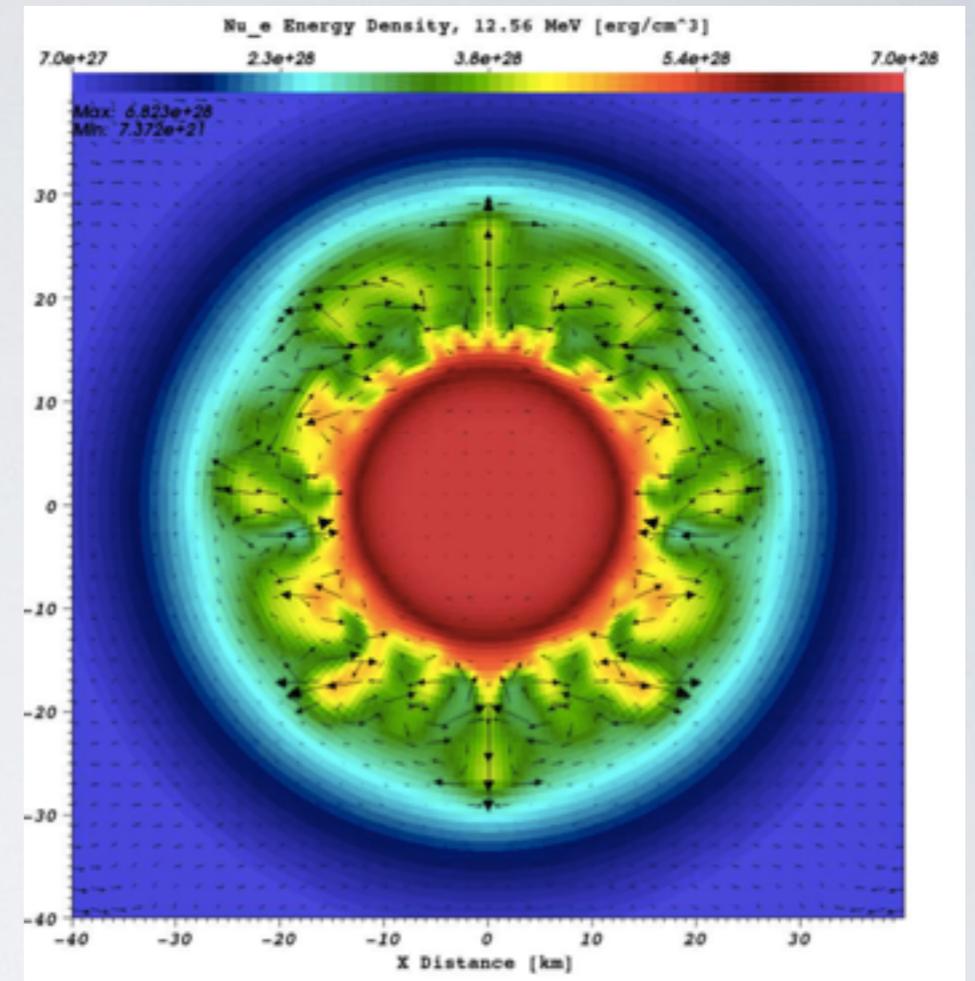
Outer material accretes supersonically, hits “wall” to produce a bounce shock.



Neutrino-matter interactions

Trapped electron neutrinos create a degenerate Fermi sea.

Stimulated absorption in reactions due to Pauli blocking (stimulated *emission* for photons).

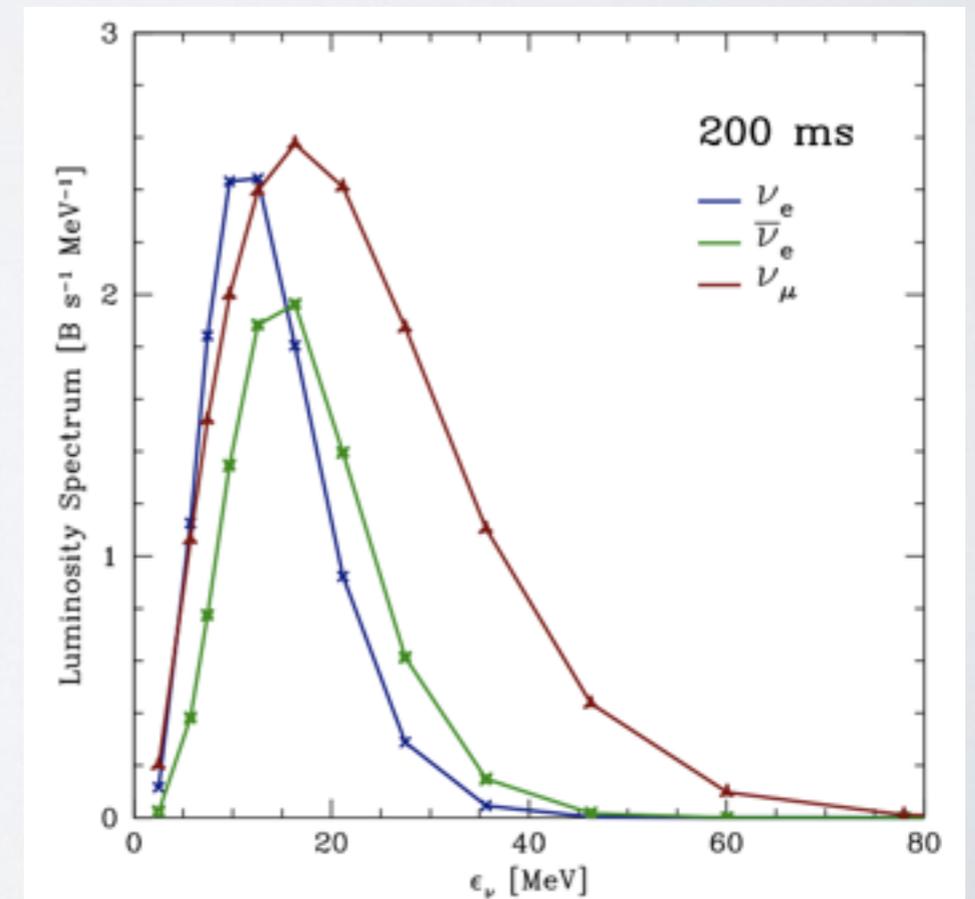


Neutrino Types:

Electron	ν_e	Anti-Electron	$\bar{\nu}_e$
Muon	ν_μ	Anti-Muon	$\bar{\nu}_\mu$
Tau	ν_τ	Anti-Tau	$\bar{\nu}_\tau$

Typical Energies:

Beta Decay Neutrinos	$\leq 0.5 \text{ MeV}$
Solar Neutrinos	$\sim 0.1 - 17 \text{ MeV}$
Supernova Neutrinos	$\sim 5 - 40 \text{ MeV}$



Neutrino-matter interactions

Neutrino-matter cross sections are very low.

Thomson electron scattering cross section: $\sim 10^{-24} \text{ cm}^2$

Weak interaction cross sections for supernova energies: $\sim 10^{-42} \text{ cm}^2$

Mean free path $\sim 100 \left(\frac{10^{15} \text{ g cm}^{-3}}{\rho} \right) \left(\frac{10 \text{ MeV}}{\epsilon_\nu} \right)^2 \text{ cm}$

Significant interactions:

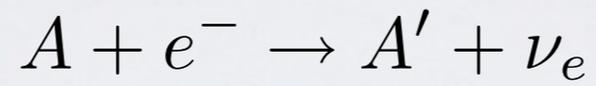
Reaction	Neutrino Type	Cross Section σ [$\times 10^{-42} \text{ cm}^2$]
$\nu_e + n \longleftrightarrow e^- + p$	Electron	$\sim 8 \left(\frac{\epsilon_\nu}{10 \text{ MeV}} \right)^2$
$\bar{\nu}_e + p \longleftrightarrow n + e^+$	Anti-electron	$\sim 7 \left(\frac{\epsilon_\nu}{10 \text{ MeV}} \right)^2$
$\nu_i + p \longrightarrow \nu_i + p$	All species	$\sim 1.7 \left(\frac{\epsilon_\nu}{10 \text{ MeV}} \right)^2$
$\nu_i + n \longrightarrow \nu_i + n$	All species	$\sim 2.0 \left(\frac{\epsilon_\nu}{10 \text{ MeV}} \right)^2$
$\nu_i + A \longrightarrow \nu_i + A$	All species	$\sim 1.2 \left(\frac{\epsilon_\nu}{10 \text{ MeV}} \right)^2$
$\nu_i + e^- \longrightarrow \nu_i + e^-$	All species	$\nu_e : \sim 5 \left(\frac{\epsilon_\nu}{10 \text{ MeV}} \right) \left(\frac{T + \mu_e/4}{10 \text{ MeV}} \right)$

Stalling of the shock

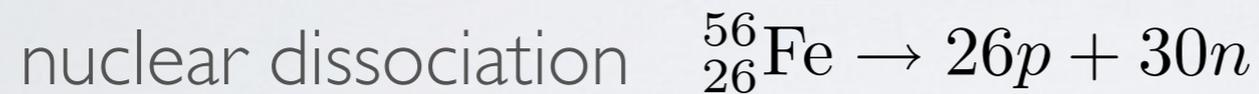
Shock wave propagates into heavy nuclei with high specific heat.

Shock energy goes into:

electron capture



nuclear dissociation



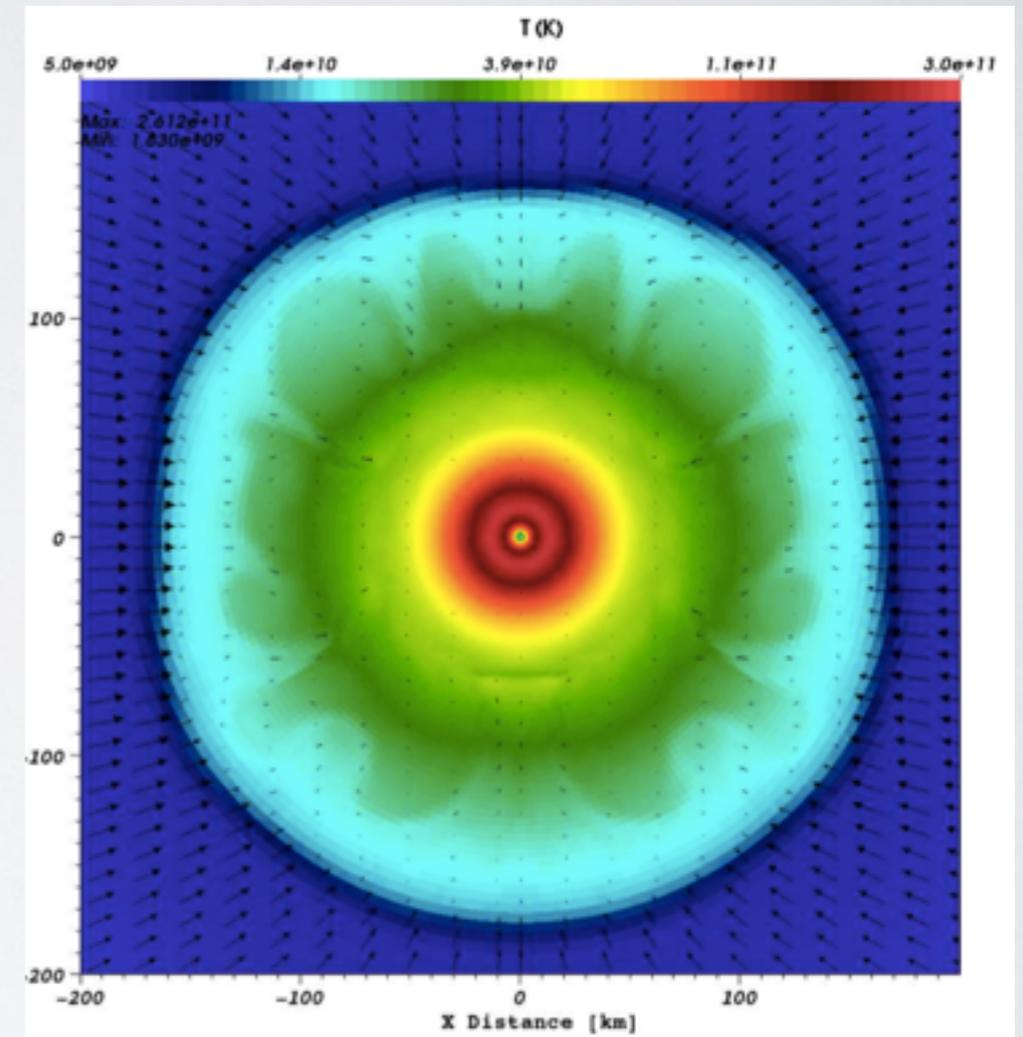
Post-shock pressure falls

Ram pressure $p_{\text{ram}} = \rho v^2$ may be more than post shock pressure

Prompt explosion will only work if ram pressure drops.

Requires a steep density gradient. $8 - 10M_{\odot}$

Otherwise $p_{\text{ram}} > p_{\text{post-shock}}$ implies that the shock stalls.



Reviving the shock

Need to either increase increase $p_{\text{post-shock}}$, decrease p_{ram} , or both by:

Depositing additional energy behind the shock

Reviving the shock

Need to either increase $p_{\text{post-shock}}$, decrease p_{ram} , or both by:

Depositing additional energy behind the shock

Changing the nuclear equation of state (no, massive neutron stars)

Reviving the shock

Need to either increase $p_{\text{post-shock}}$, decrease p_{ram} , or both by:

Depositing additional energy behind the shock

Changing the nuclear equation of state (no, massive neutron stars)

Using radiation pressure (no, $L \sim 10^{53} \text{ erg s}^{-1} \ll L_{\text{edd}} \sim 10^{55} \text{ erg s}^{-1}$)

Reviving the shock

Need to either increase $p_{\text{post-shock}}$, decrease p_{ram} , or both by:

Depositing additional energy behind the shock

Changing the nuclear equation of state (no, massive neutron stars)

Using radiation pressure (no, $L \sim 10^{53} \text{ erg s}^{-1} \ll L_{\text{edd}} \sim 10^{55} \text{ erg s}^{-1}$)

Using progenitor models with steeper density profiles (no).

Reviving the shock

Need to either increase increase $p_{\text{post-shock}}$, decrease p_{ram} , or both by:

Depositing additional energy behind the shock

Possible energy sources:

Nuclear burning (O into Fe).

Releases $\sim 5 \times 10^{17} \text{ erg g}^{-1} \ll 10^{19} \text{ erg g}^{-1} \sim U_{\text{grav,be}}(r = 200 \text{ km})$

For 10^{51} erg need to burn $1M_{\odot}$ of oxygen explosively, implies need to already have enormous blast wave at large radii. See Kushnir 2015.

Annihilation: $\nu_e + \bar{\nu}_e \rightarrow e^+ + e^- \rightarrow 2\gamma$

Most efficient where neutrino-cooling is severe

Delayed neutrino heating: the neutrino mechanism.

Neutrino Transport

Simplified neutrino radiation transport, neglecting scattering:

$$\frac{\partial I_\epsilon(r)}{\partial r} = \kappa_\epsilon \rho [B_\epsilon(T) - I_\epsilon(r)]$$

$B_\epsilon(T)$ is the neutrino blackbody function.

Integrated over ϵ_ν :
$$\frac{\partial F(r)}{\partial r} = \bar{\kappa} \rho [acT^4 - F(r)]$$

$\bar{\kappa} \rho acT^4$ is emission from matter (cooling).

$\bar{\kappa} \rho F(r)$ is absorption by matter (heating).

Neutrino Cooling

Neutrino cooling dominated by URCA processes:

Electron capture: $p + e^- \rightarrow \nu_e + n$

Positron capture: $n + e^+ \rightarrow \bar{\nu}_e + p$

Neutrino Cooling

Neutrino cooling dominated by URCA processes:

Electron capture: $p + e^- \rightarrow \nu_e + n$

Positron capture: $n + e^+ \rightarrow \bar{\nu}_e + p$

Neutrino energy loss rate per gram: $\bar{\kappa} a c T^4$

$$\bar{\kappa} \propto \epsilon_\nu^2 \Rightarrow \bar{\kappa} \propto T^2$$

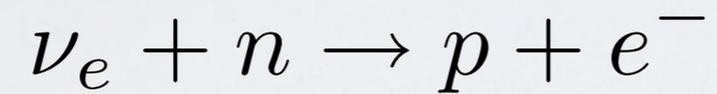
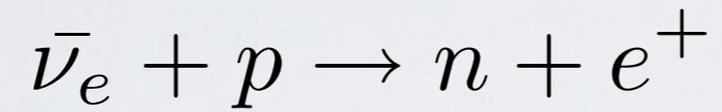
$$\text{Cooling rate per gram: } \mathcal{C} \approx 1.4 \times 10^{20} \left(\frac{T}{2 \text{ MeV}} \right)^6 \text{ erg g}^{-1} \text{ s}^{-1}$$

Cooling dominates at $R < 80$ km where matter is hot.

Neutrino Heating

Neutrino heating rate per gram: $\bar{\kappa}F_\nu$

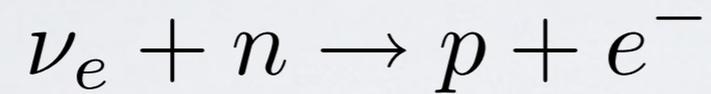
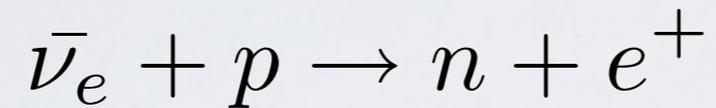
Neutrino capture, inverses of URCA processes:



Neutrino Heating

Neutrino heating rate per gram: $\bar{\kappa}F_\nu$

Neutrino capture, inverses of URCA processes:



$$\bar{\kappa} \propto \epsilon_\nu^2 \Rightarrow \bar{\kappa} \propto T_{\nu_e}^2$$

Heating rate per gram: $\mathcal{H} \propto T_{\nu_e}^2 \frac{L_\nu}{4\pi r^2}$

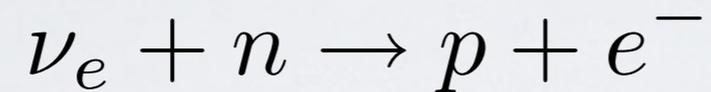
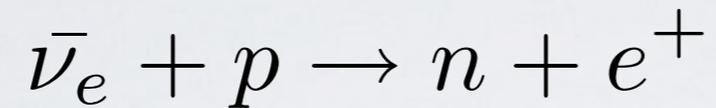
$$\mathcal{H} \approx 1.5 \times 10^{20} L_{\nu_e} \left(\frac{100 \text{ km}}{r} \right)^2 \left(\frac{T_{\nu_e}}{4 \text{ MeV}} \right)^2 \text{ erg g}^{-1} \text{ s}^{-1}$$

Characteristic temperature of the neutrino spectrum, does not decrease with distance.

Neutrino Heating

Neutrino heating rate per gram: $\bar{\kappa}F_\nu$

Neutrino capture, inverses of URCA processes:



$$\bar{\kappa} \propto \epsilon_\nu^2 \Rightarrow \bar{\kappa} \propto T_{\nu_e}^2$$

Heating rate per gram: $\mathcal{H} \propto T_{\nu_e}^2 \frac{L_\nu}{4\pi r^2}$

$$\mathcal{H} \approx 1.5 \times 10^{20} L_{\nu_e} \left(\frac{100 \text{ km}}{r} \right)^2 \left(\frac{T_{\nu_e}}{4 \text{ MeV}} \right)^2 \text{ erg g}^{-1} \text{ s}^{-1}$$

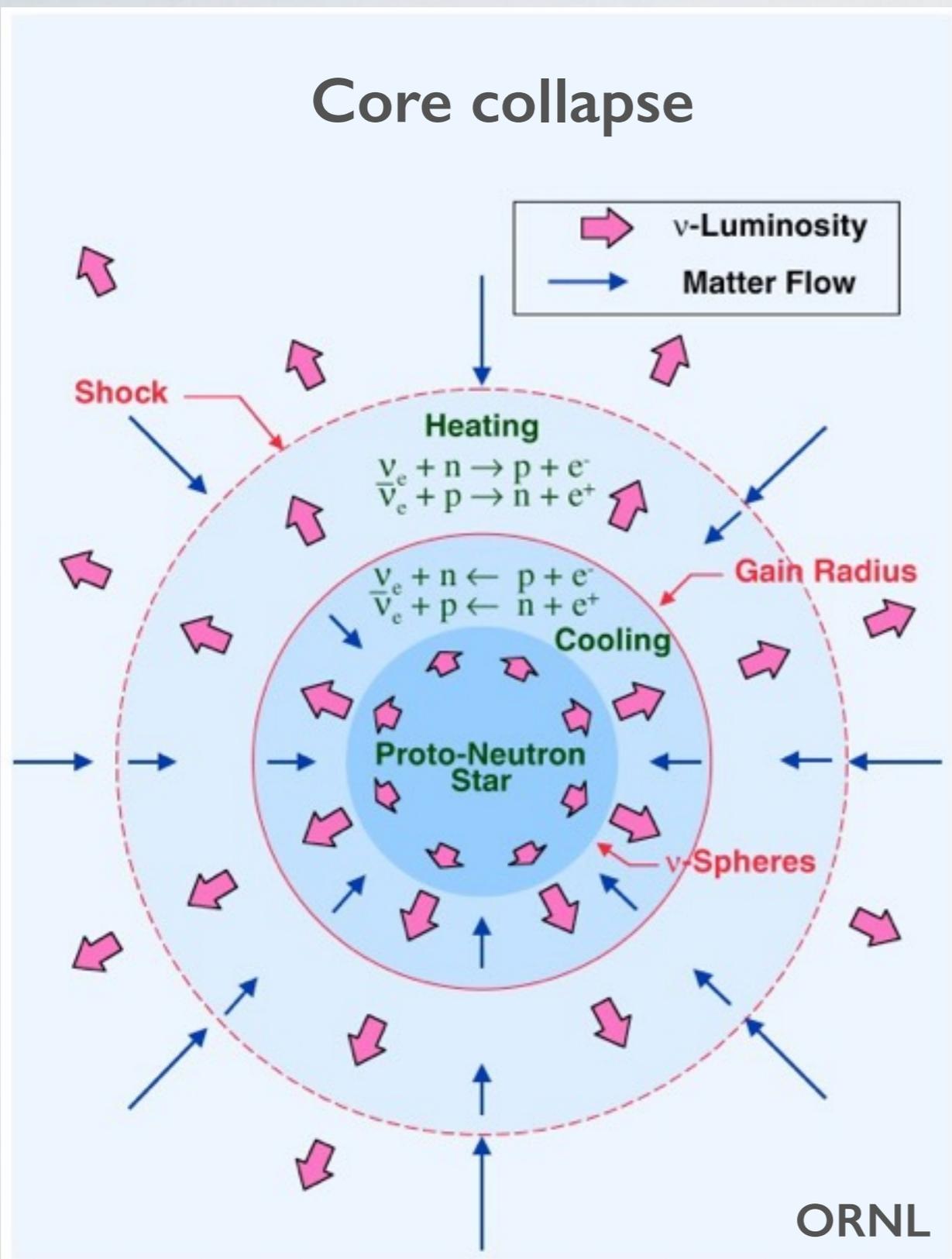
Characteristic temperature of the neutrino spectrum, does not decrease with distance.

Gain region: Net neutrino heating

$$\mathcal{H} - \mathcal{C} > 0 \Rightarrow 80 \text{ km} - R_{\text{shock}}$$

Gain Region

Core collapse



Gain region: Net neutrino heating

$$\mathcal{H} - \mathcal{C} > 0 \Rightarrow 80 \text{ km} - R_{\text{shock}}$$

Does it work?

Net neutrino heating adds $\sim 10^{20} \text{ erg g}^{-1} \text{ s}^{-1}$ to the gain region.

Matter gains sufficient energy to become unbound in $\sim 100 \text{ ms}$.

Timescales and Quantities

Quantity	Definition	Description
Residence Time	τ_{res}	Mean dwell time in the gain region.
Heating Time	$\tau_{\text{q}} \equiv \frac{\int_{\text{gain}} U_{\text{int}}}{\int_{\text{gain}} (\mathcal{H} - \mathcal{C})}$	Characteristic timescale to heat gain region.
Heating Power	$\mathcal{P} \equiv \tau L_{\nu}$	Net neutrino energy deposition rate.

Timescales and Quantities

Quantity	Definition	Description
Residence Time	τ_{res}	Mean dwell time in the gain region.
Heating Time	$\tau_{\text{q}} \equiv \frac{\int_{\text{gain}} U_{\text{int}}}{\int_{\text{gain}} (\mathcal{H} - \mathcal{C})}$	Characteristic timescale to heat gain region.
Heating Power	$\mathcal{P} \equiv \tau L_{\nu}$	Net neutrino energy deposition rate.

Important dimensionless ratio: $\frac{\tau_{\text{res}}}{\tau_{\text{q}}}$

$\frac{\tau_{\text{res}}}{\tau_{\text{q}}} > 1$ implies net energy and pressure added to gain region.

Timescales and Quantities

Quantity	Definition	Description
Residence Time	τ_{res}	Mean dwell time in the gain region.
Heating Time	$\tau_q \equiv \frac{\int_{\text{gain}} U_{\text{int}}}{\int_{\text{gain}} (\mathcal{H} - \mathcal{C})}$	Characteristic timescale to heat gain region.
Heating Power	$\mathcal{P} \equiv \tau L_\nu$	Net neutrino energy deposition rate.

Important dimensionless ratio: $\frac{\tau_{\text{res}}}{\tau_q}$

$\frac{\tau_{\text{res}}}{\tau_q} > 1$ implies net energy and pressure added to gain region.

If sustained, shock is revived.

$$L_\nu - \dot{\mathcal{M}} \quad \text{relation} \quad \frac{\tau_{\text{res}}}{\tau_q} \sim \frac{\mathcal{M}_{\text{gain}}}{\dot{\mathcal{M}}} \frac{L_\nu}{\int_{\text{gain}} U_{\text{int}}} \sim \frac{L_\nu}{\dot{\mathcal{M}}} \frac{m_p}{k_B T_{\text{gain}}}$$

Timescales and Quantities

Quantity	Definition	Description
Residence Time	τ_{res}	Mean dwell time in the gain region.
Heating Time	$\tau_{\text{q}} \equiv \frac{\int_{\text{gain}} U_{\text{int}}}{\int_{\text{gain}} (\mathcal{H} - \mathcal{C})}$	Characteristic timescale to heat gain region.
Heating Power	$\mathcal{P} \equiv \tau L_{\nu}$	Net neutrino energy deposition rate.

Important dimensionless ratio: $\frac{\tau_{\text{res}}}{\tau_{\text{q}}}$

$\frac{\tau_{\text{res}}}{\tau_{\text{q}}} > 1$ implies net energy and pressure added to gain region.

If sustained, shock is revived.

$$L_{\nu} - \dot{\mathcal{M}} \quad \text{relation} \quad \frac{\tau_{\text{res}}}{\tau_{\text{q}}} \sim \frac{\mathcal{M}_{\text{gain}}}{\dot{\mathcal{M}}} \frac{L_{\nu}}{\int_{\text{gain}} U_{\text{int}}} \sim \frac{L_{\nu}}{\dot{\mathcal{M}}} \frac{m_p}{k_B T_{\text{gain}}}$$

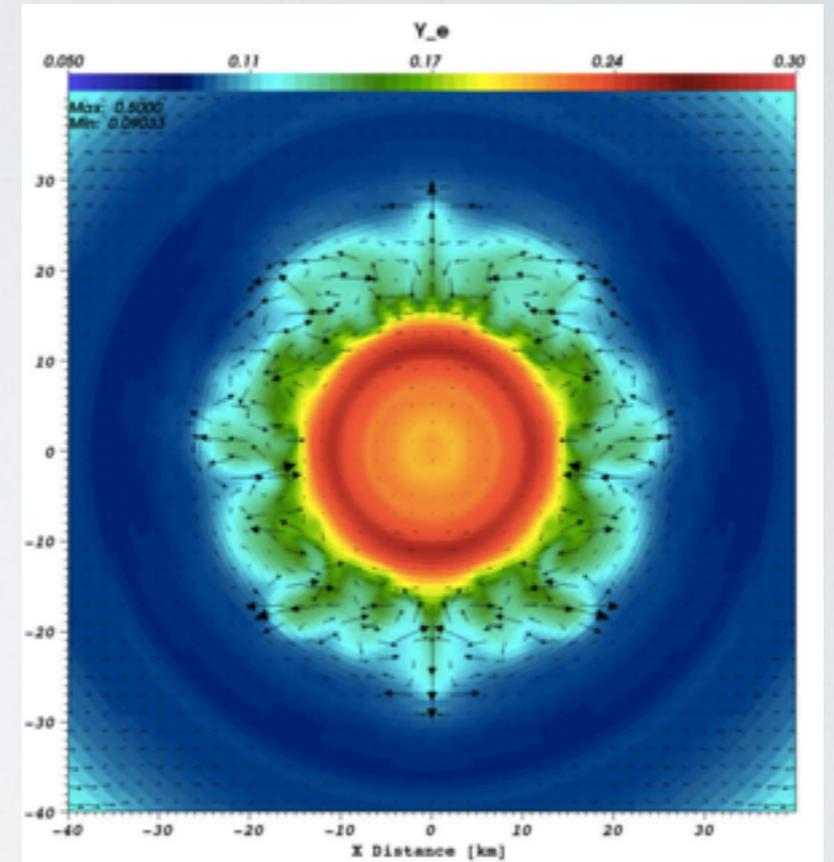
To make neutrino mechanism work, need to increase τ_{res} and/or L_{ν} .

Convection and GR: Increasing $\mathcal{P} \equiv \tau L_\nu$

Core convection can dredge up trapped neutrinos, potentially increasing the neutrino luminosities.

Not seen in high-res multi-D simulations.

Outer core convection only.



Convection and GR: Increasing $\mathcal{P} \equiv \tau L_\nu$

Core convection can dredge up trapped neutrinos, potentially increasing the neutrino luminosities.

Not seen in high-res multi-D simulations.

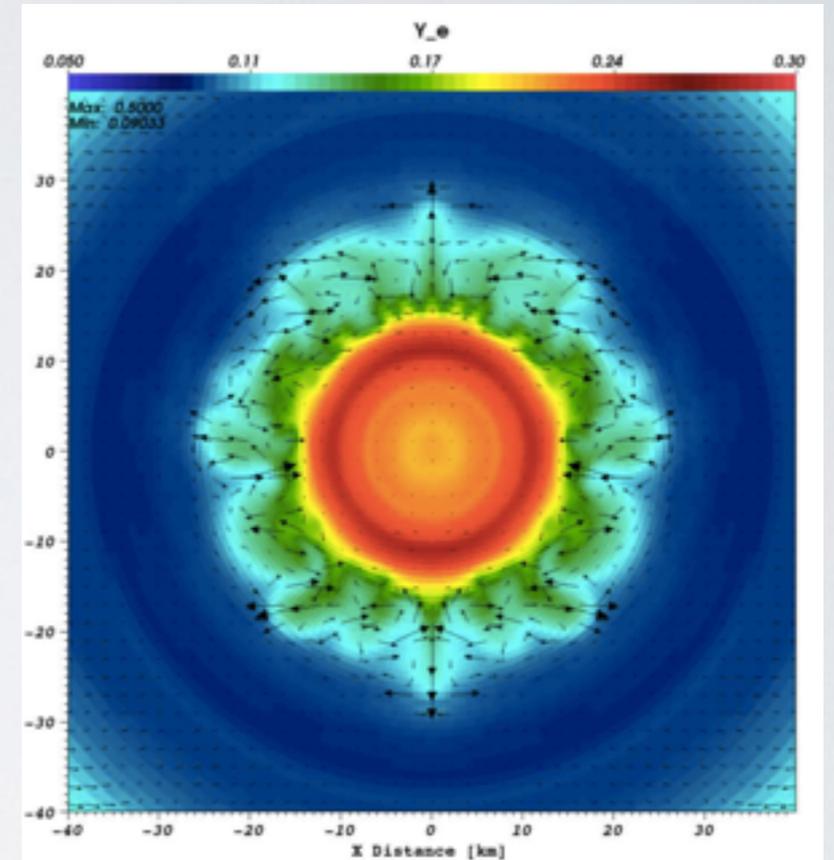
Outer core convection only.

General Relativity

Relativity tends to make the core more compact.

Hotter neutrino-spheres imply higher cross-sections which imply more heating.
But neutrinos are redshifted from the deeper well...

GR effects seem to give minor corrections to \mathcal{P} .



Multidimensional Effects: Increasing τ_{res}

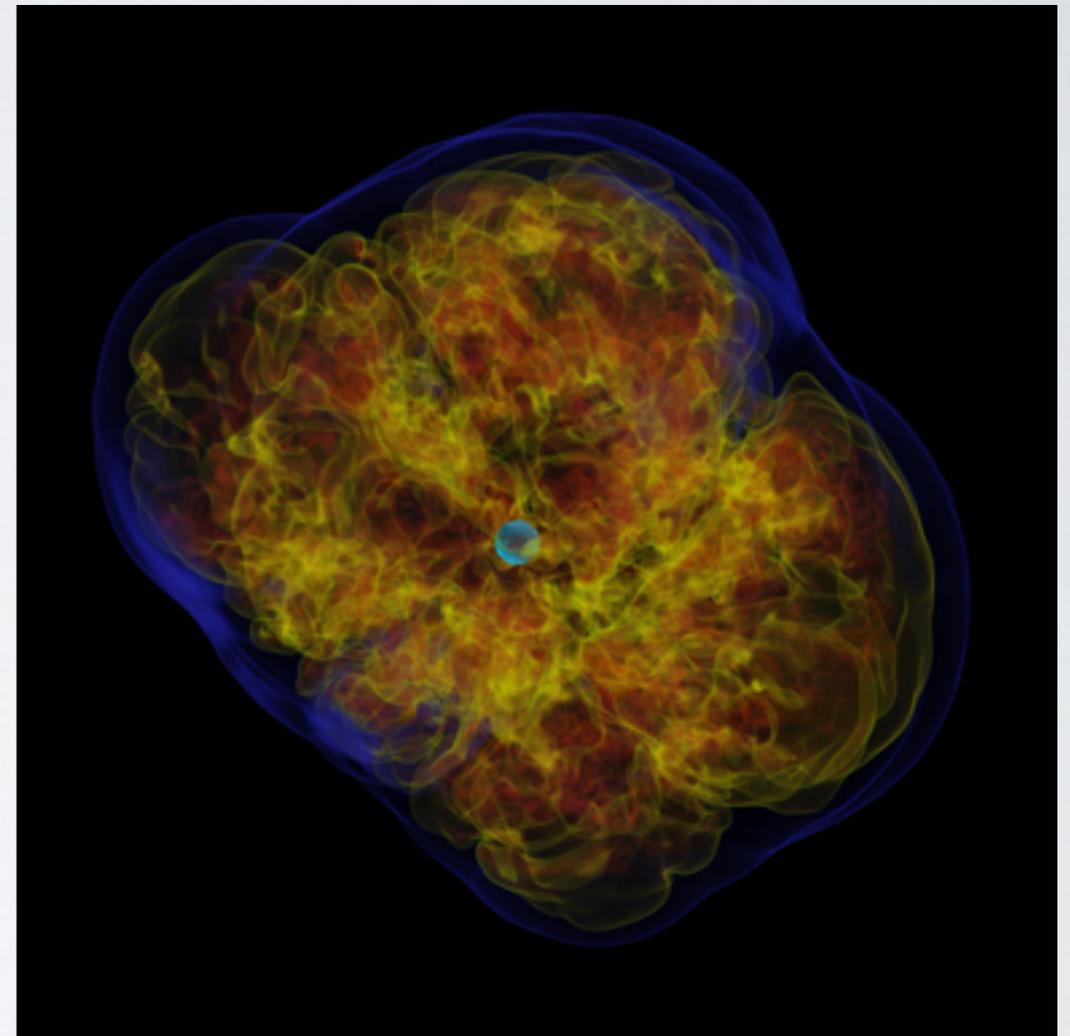
Enormous computing power: Fully 3D simulations.

Fundamentally different results in spherically symmetric, axisymmetric, 3D simulations.

Dynamics different.

Dwell times different.

Turbulence different.



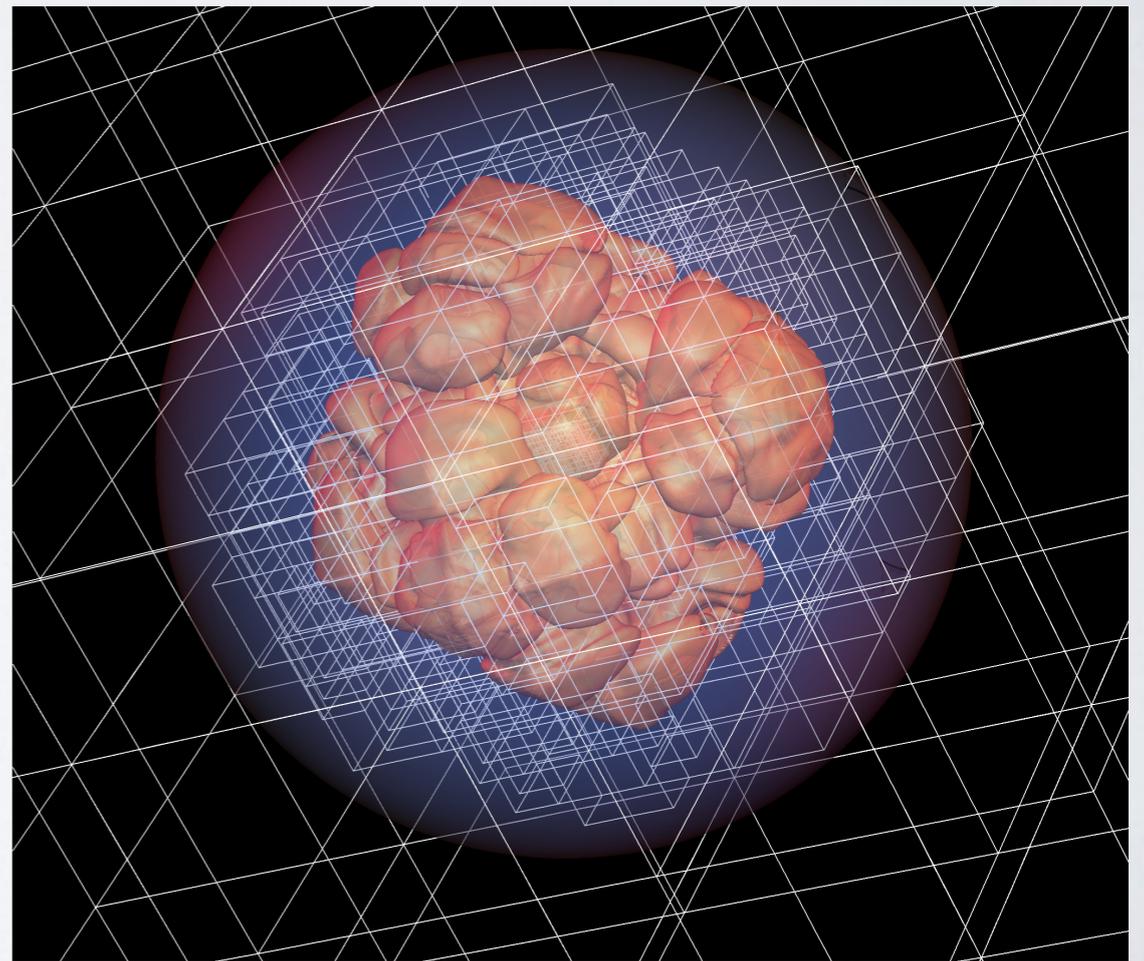
CASTRO: Compressible Astrophysics

- Multi-D radiation-hydrodynamics code
- Adaptive mesh refinement (AMR) with sub-cycling in time
- Advection: 2nd order, unsplit piecewise-linear or PPM
- Radiation: multi-group flux limited diffusion
- Gravity: Monopole or multi-grid Poisson solve
- Scales to over 200,000 cores!

Ann Almgren (LBL)

John Bell (LBL)

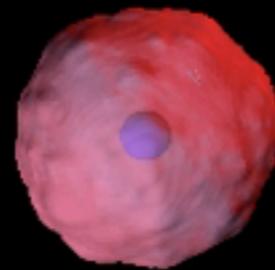
Louis Howell (LLNL)



3D AMR block structure

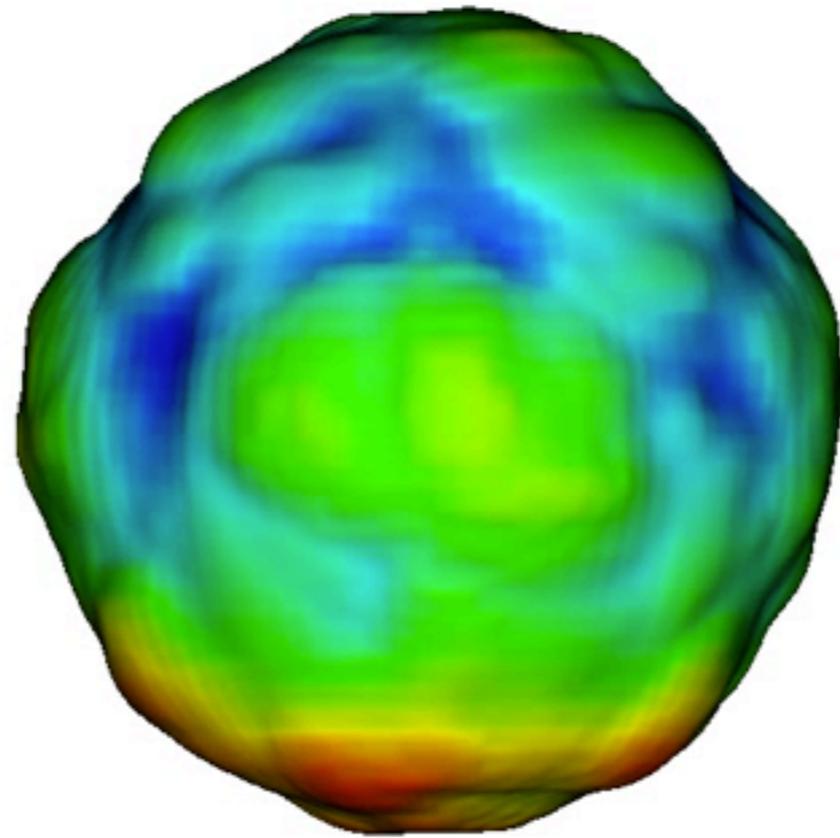
Non-Rotating Initial Model

Time=0.144418 s



2000 kilometers

150 km



Time = 0.0811 seconds

Standing Accretion Shock Instability (SASI)

Axisymmetric

$\ell = 1$ mode
is dominant

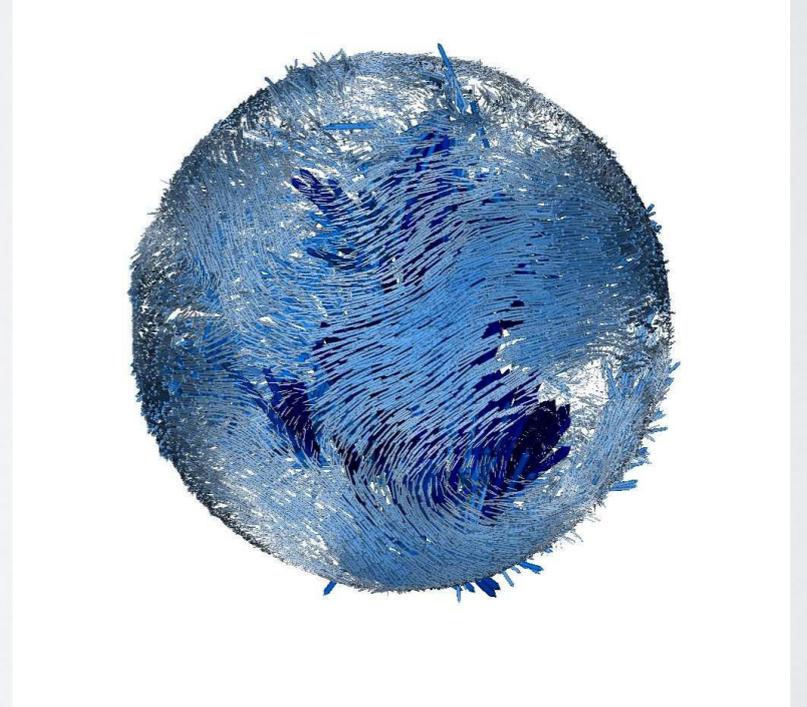
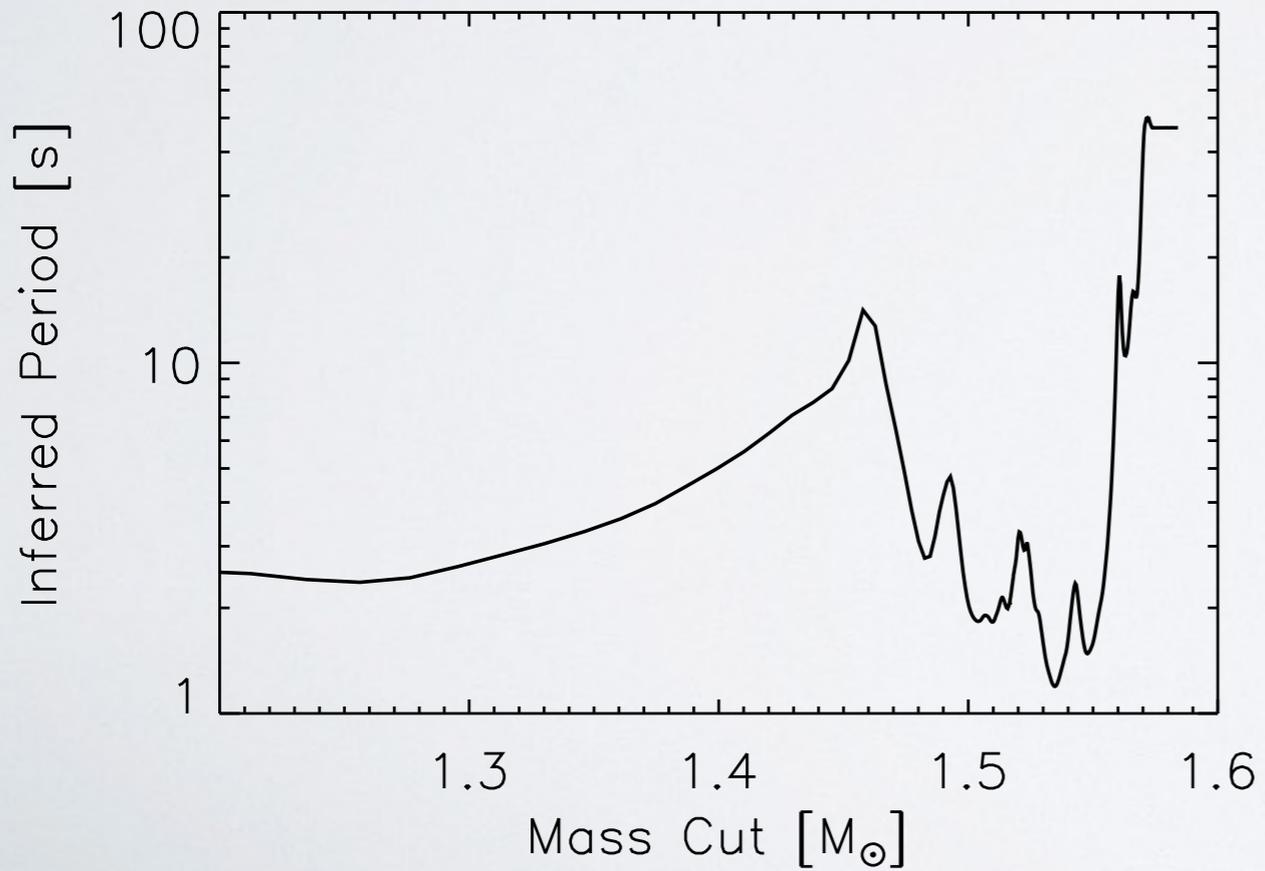
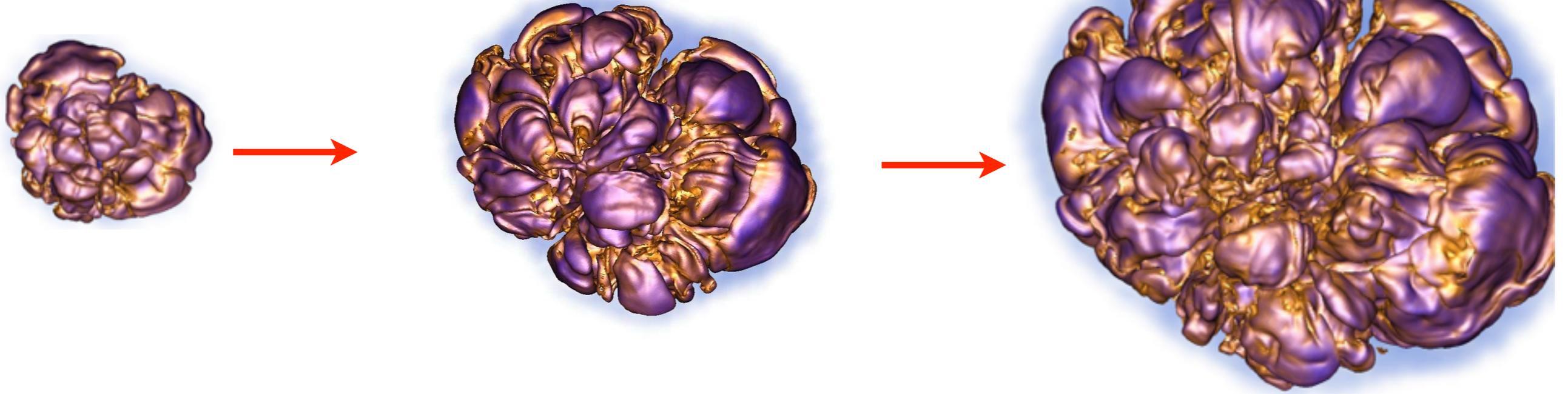
Suggested as a
fundamental
characteristic of SN
dynamics and
way to spin-up
pulsars;

Blondin & Mezzacappa 2007



2000 kilometers

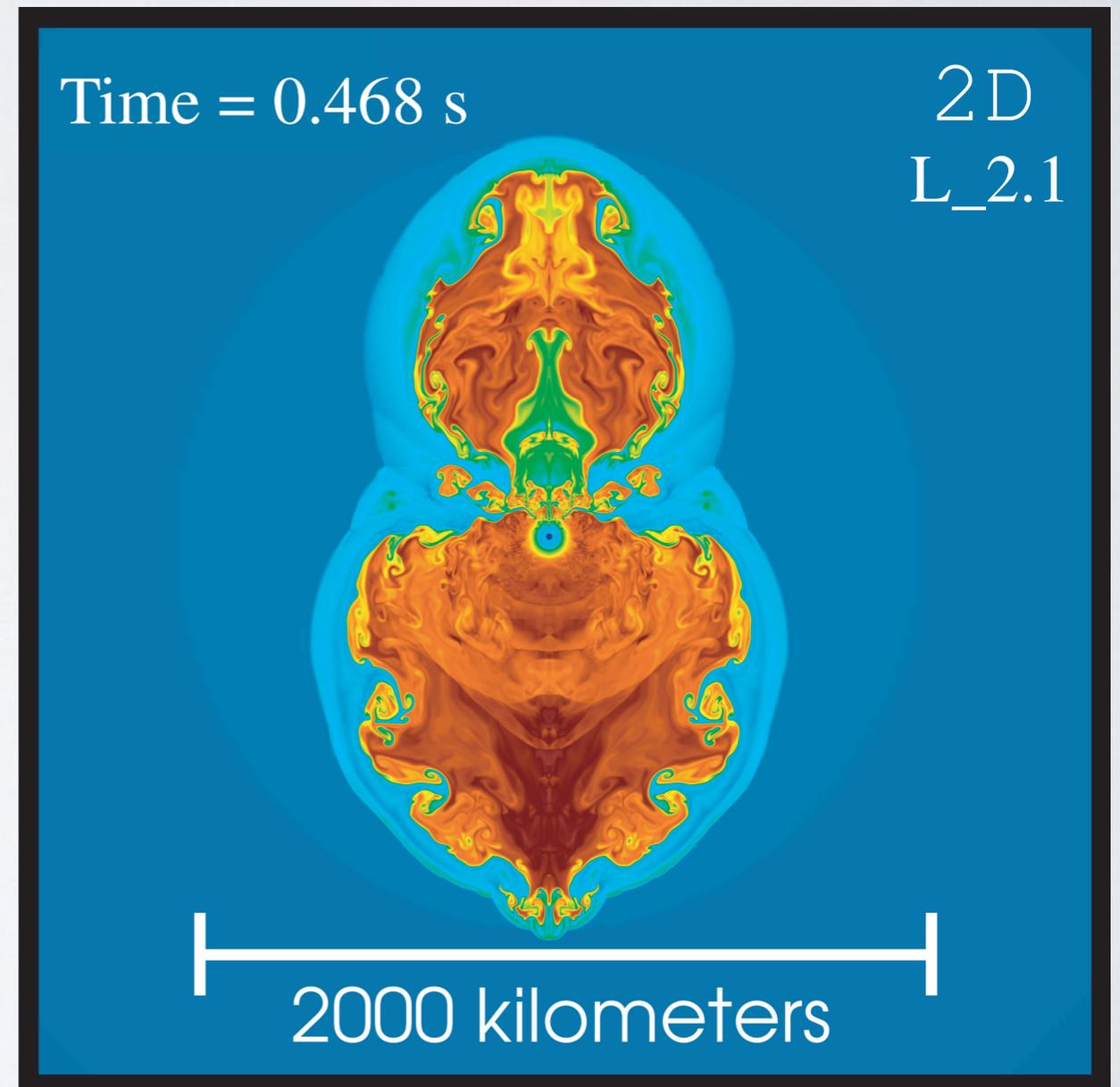
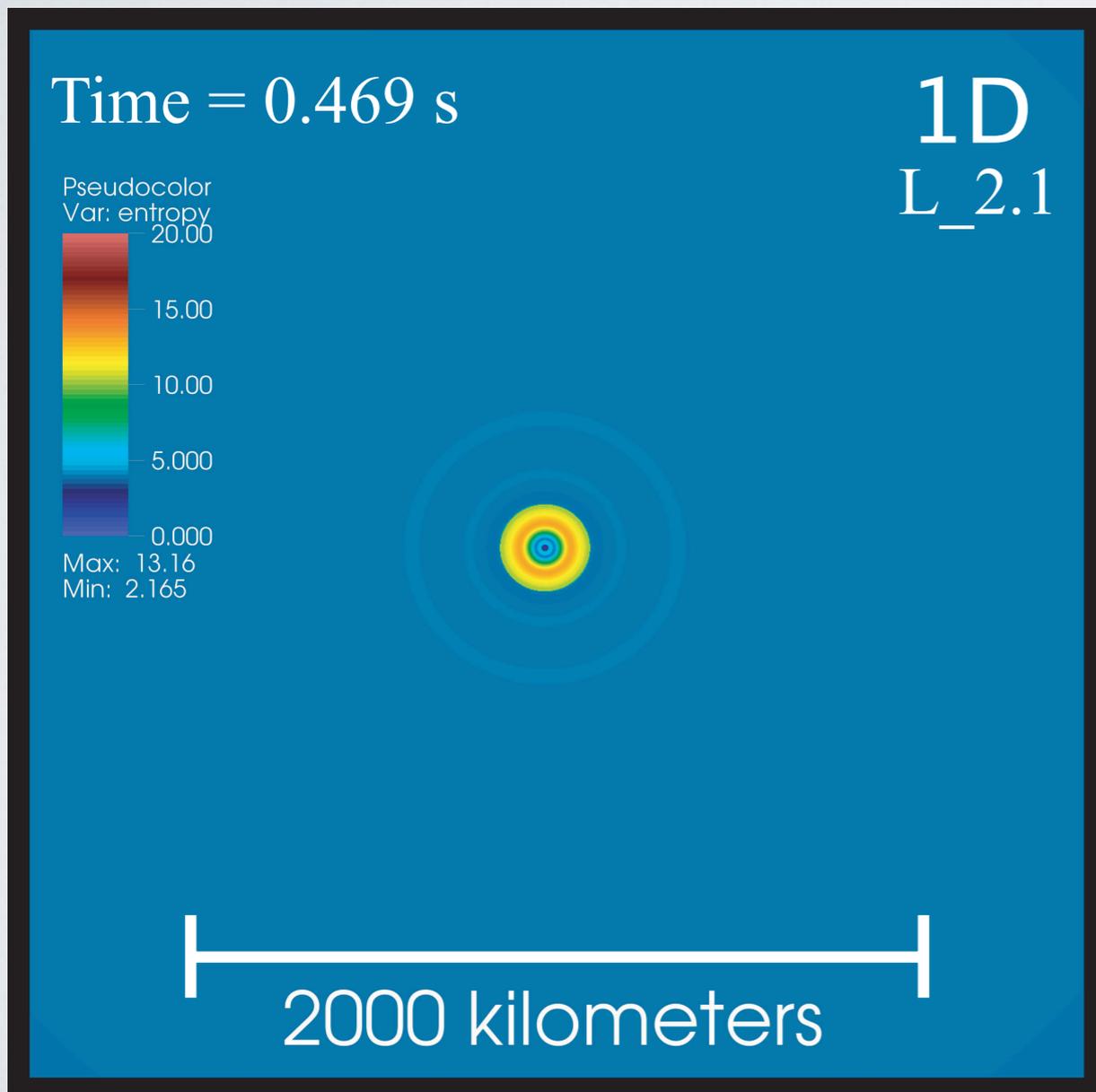
Time sequence



Dimensional Dependence

Spherically Symmetric

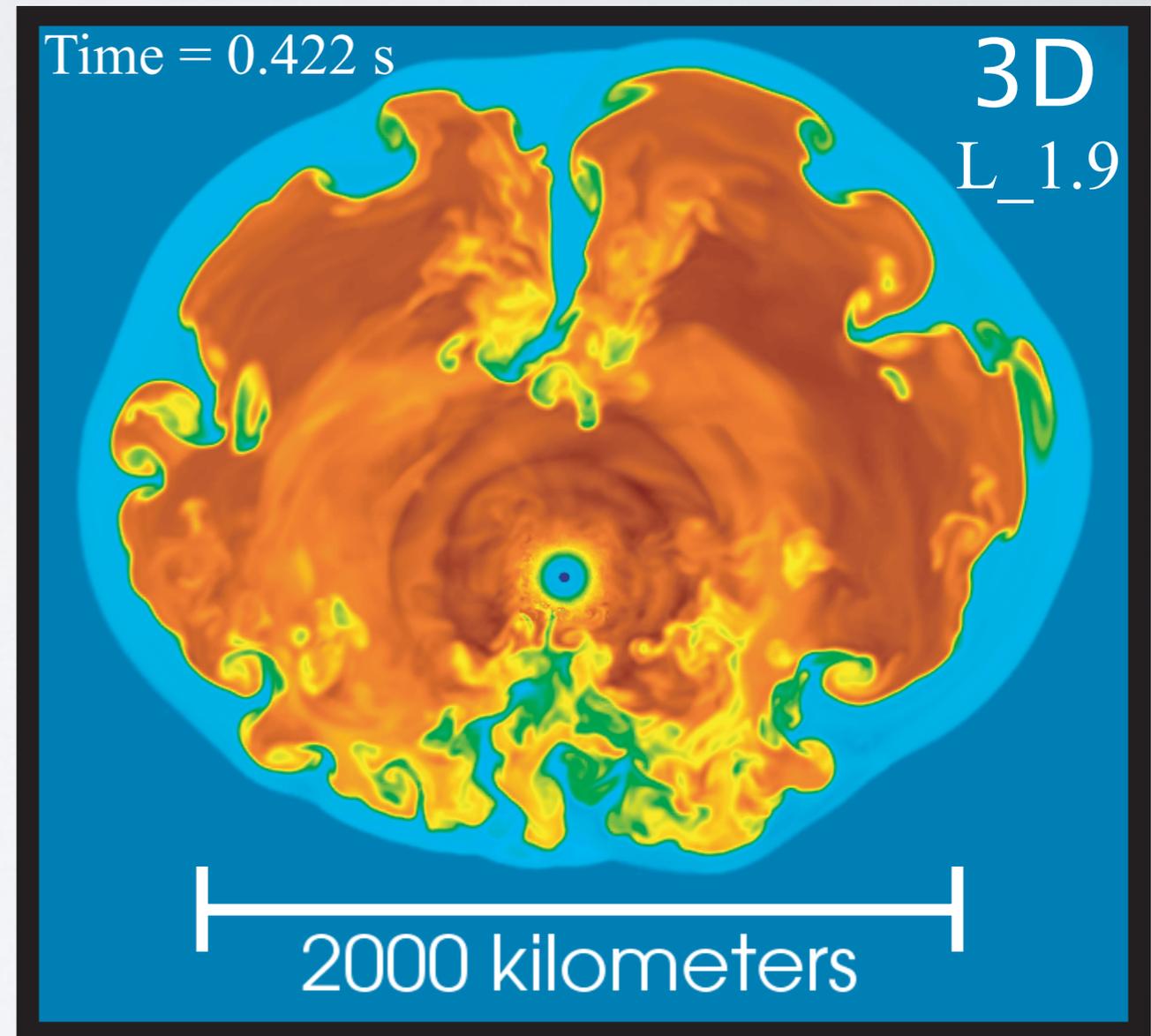
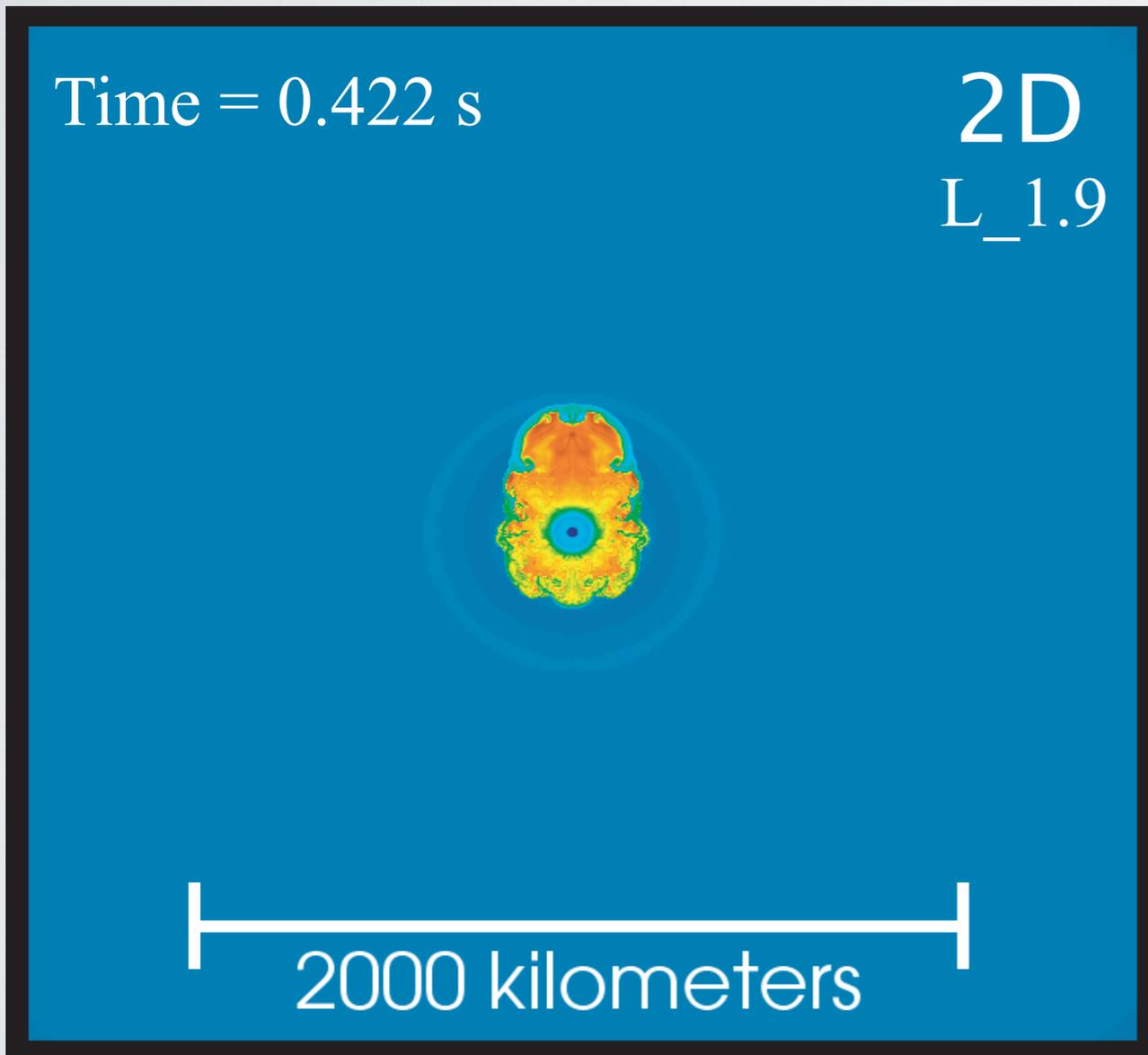
Axisymmetric

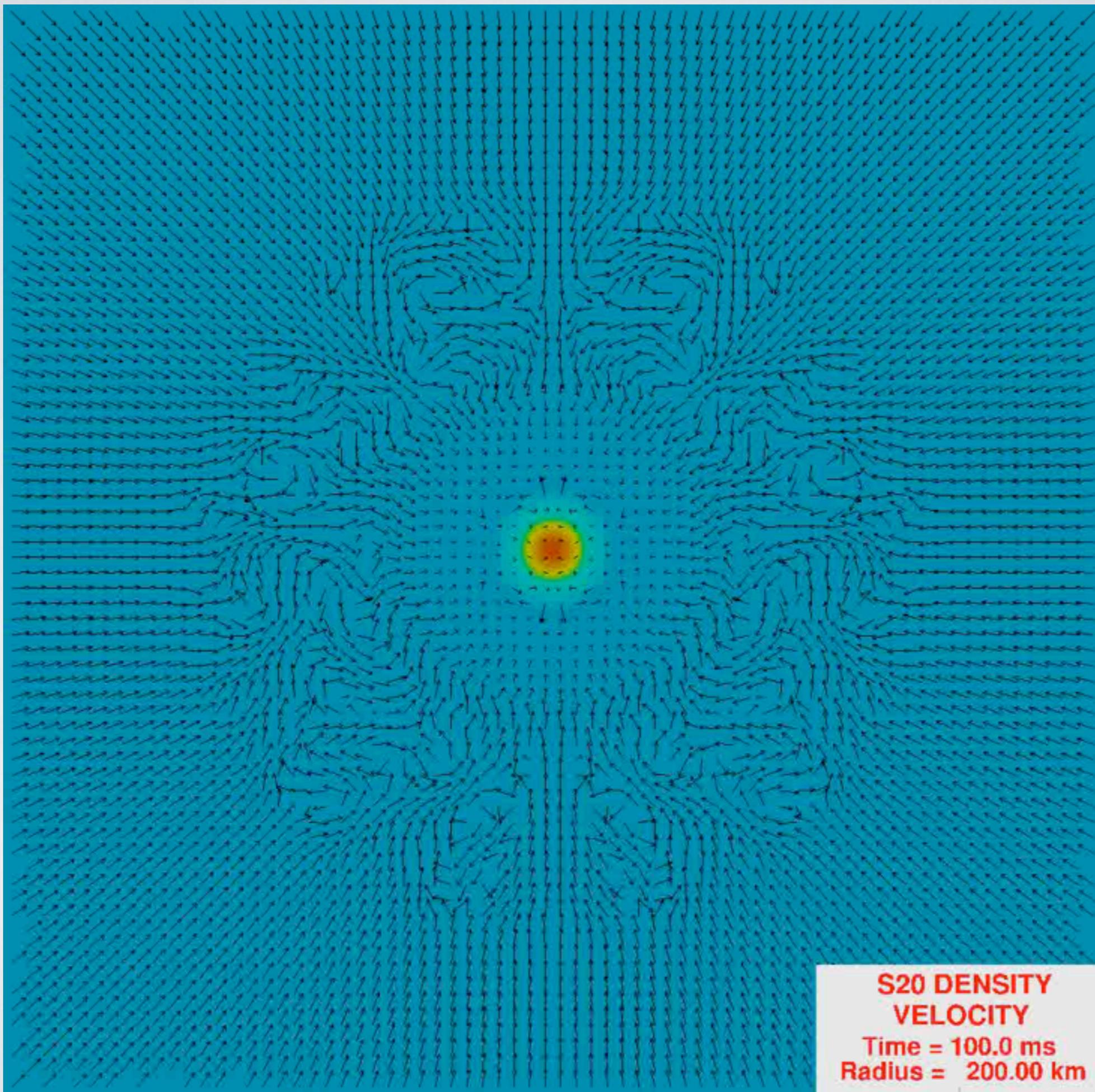


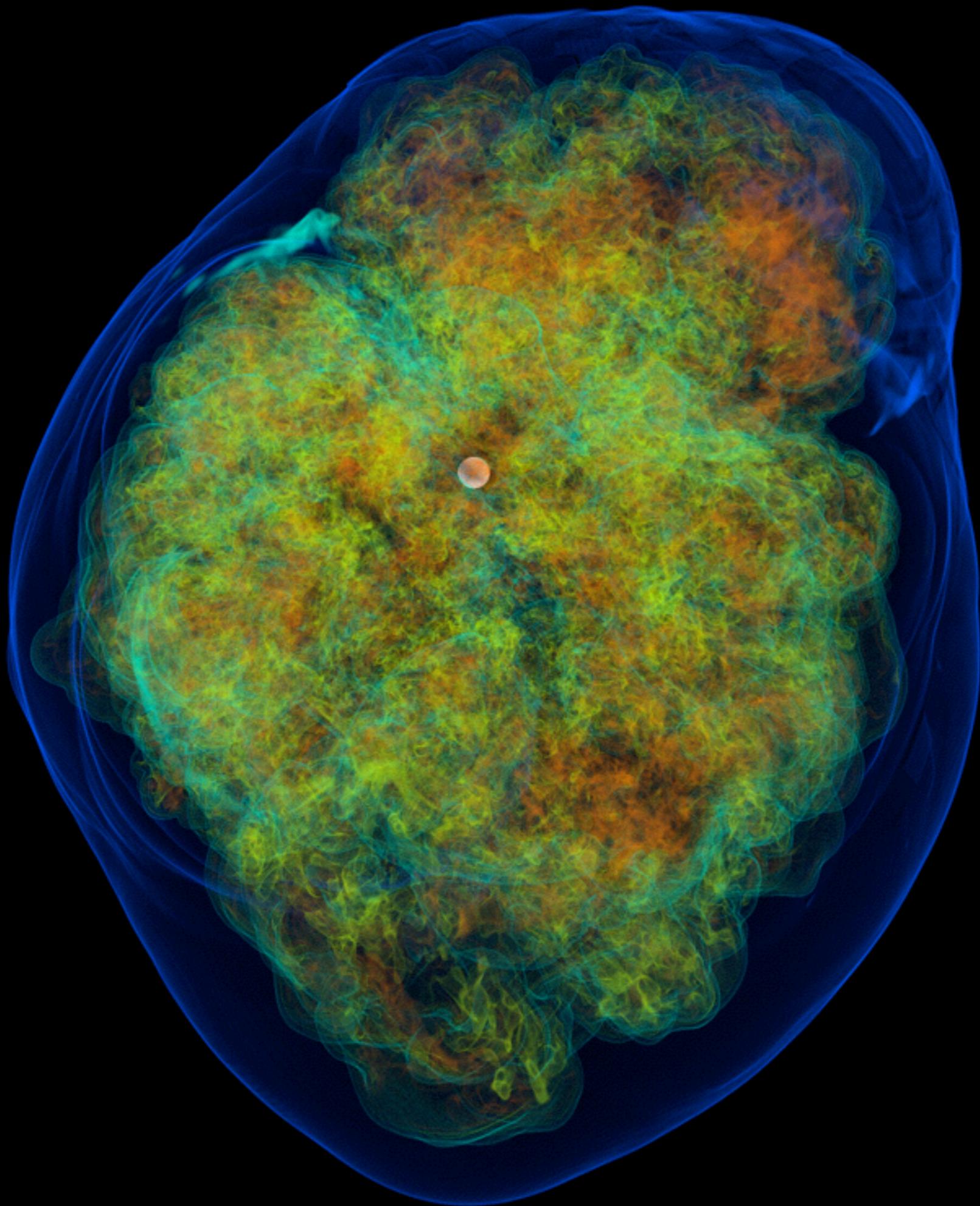
Dimensional Dependence

Axisymmetric

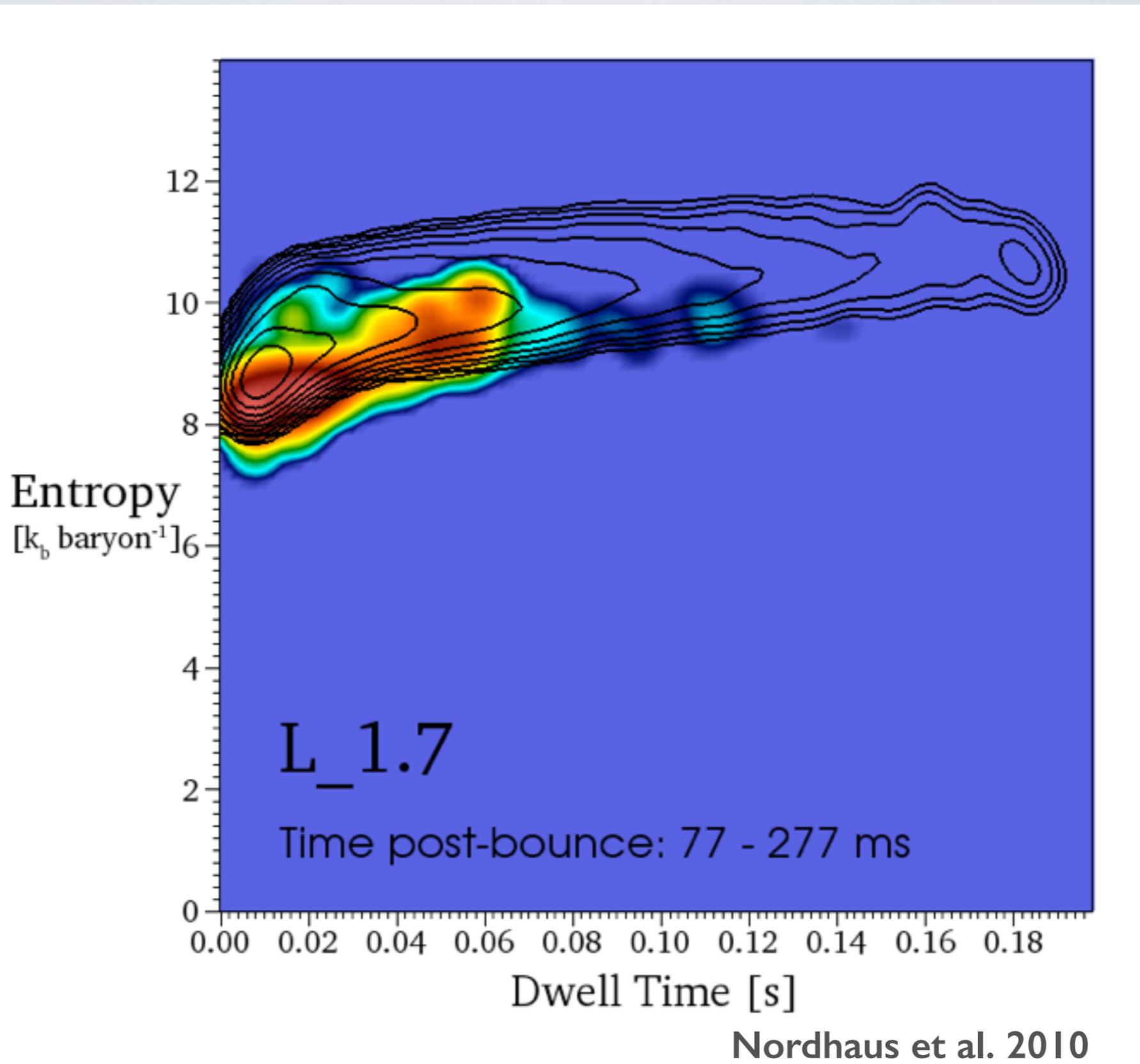
Three Dimensional



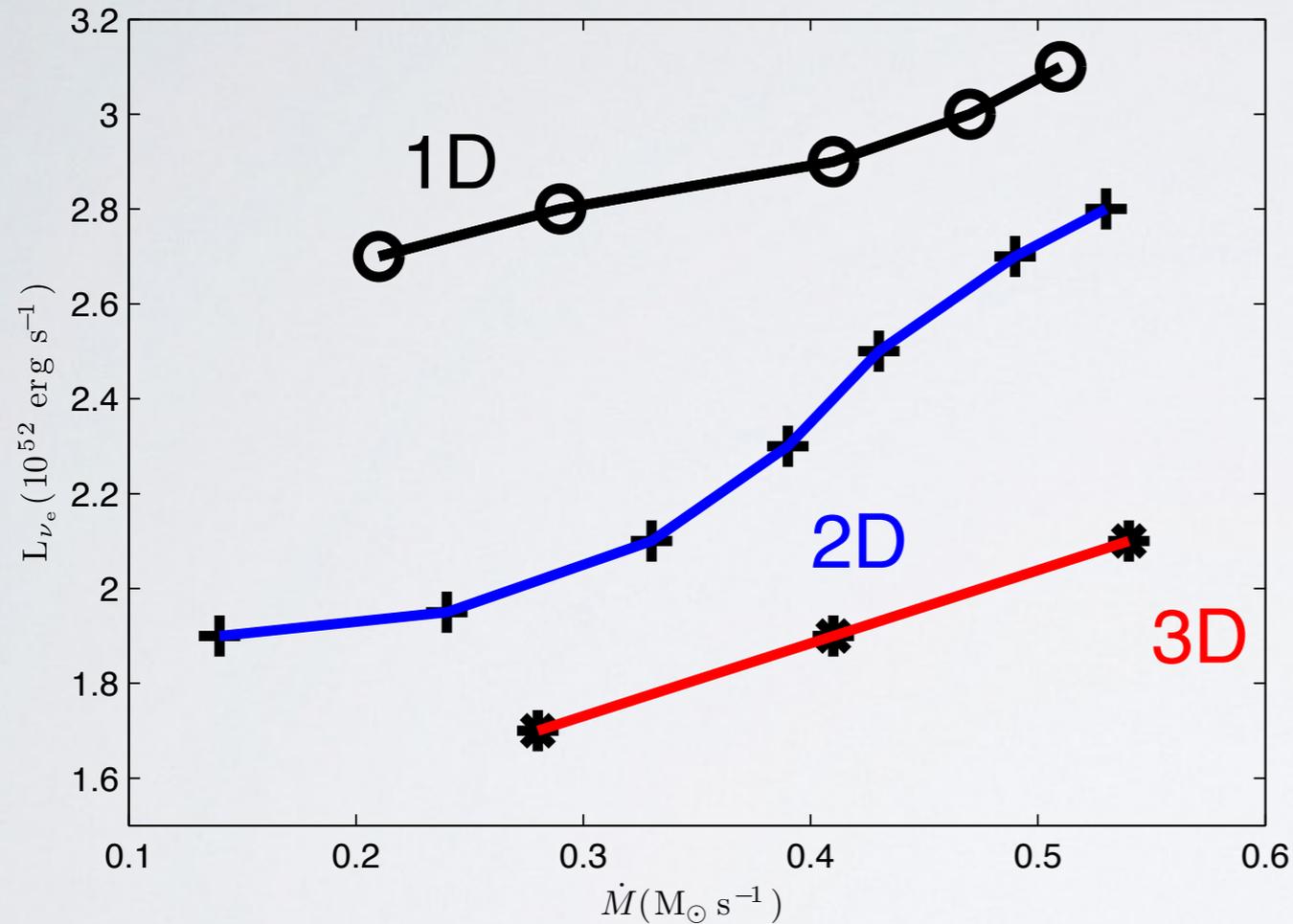




Multidimensional Effects: Increasing τ_{res}



Critical curve for explosions



3D explodes earlier

Nordhaus et al. 2010
Dolence et al. 2014
Burrows et al. 2013
Hanke et al. 2012

2D and 3D explode same time

Hanke et al. 2012
Melson et al. 2014

2D explodes earlier

Couch 2014
Couch & O'Connor 2013
Takiwaki et al. 2014

Controversial - different groups, different codes, different results.

Despite all this hope...

All 3D simulations to date:

if explosion, **under-energetic by factors of 10-100.**

Inherent limitation of the neutrino-mechanism?

Under the most extremely favorable conditions, neutrinos may accelerate the gas to energy of $5e50$ erg.

Papish, Nordhaus, Soker 2015 MNRAS

Alternate ideas.

General relativistic effects: to date mostly 2D and conformally flat.

MHD: linear winding and compression only.

Nucleosynthesis: if neutrinos fail, perhaps nuclear burning.

Kushnir 2015

Jets, winds: angular momentum present, jets ubiquitous in astrophysics.

see Jittering-Jet model/papers of Oded Papish

What questions do we aim to answer?

1. Canonical explosion energy of 1 Bethe.
2. Neutron star masses and distribution.
3. Nucleosynthetic yields and distributions.
4. Pulsar Kicks.

What questions do we aim to answer?

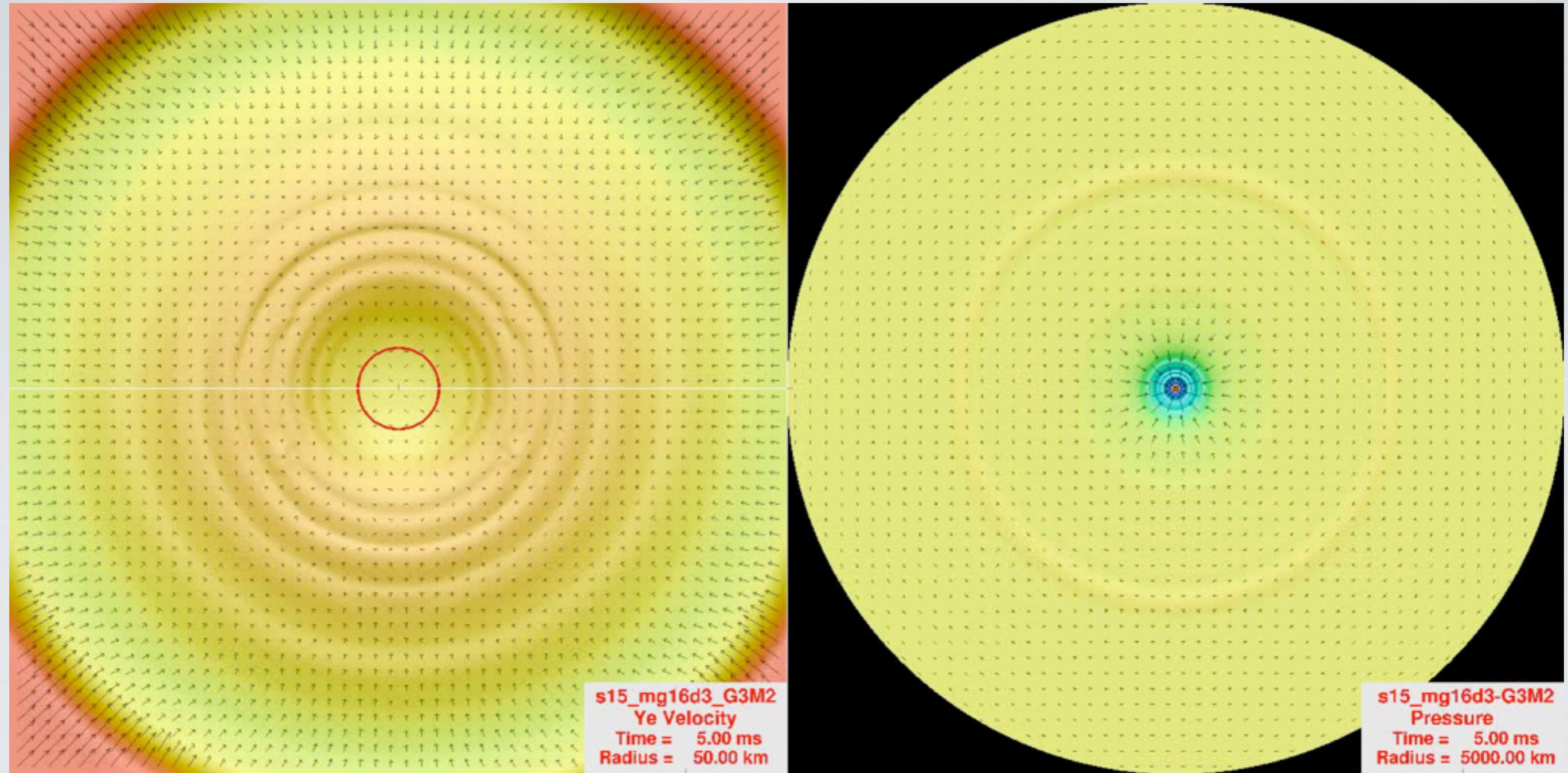
~~1. Canonical explosion energy of 1 Bethe.~~

~~2. Neutron star masses and distribution.~~

~~3. Nucleosynthetic yields and distributions.~~

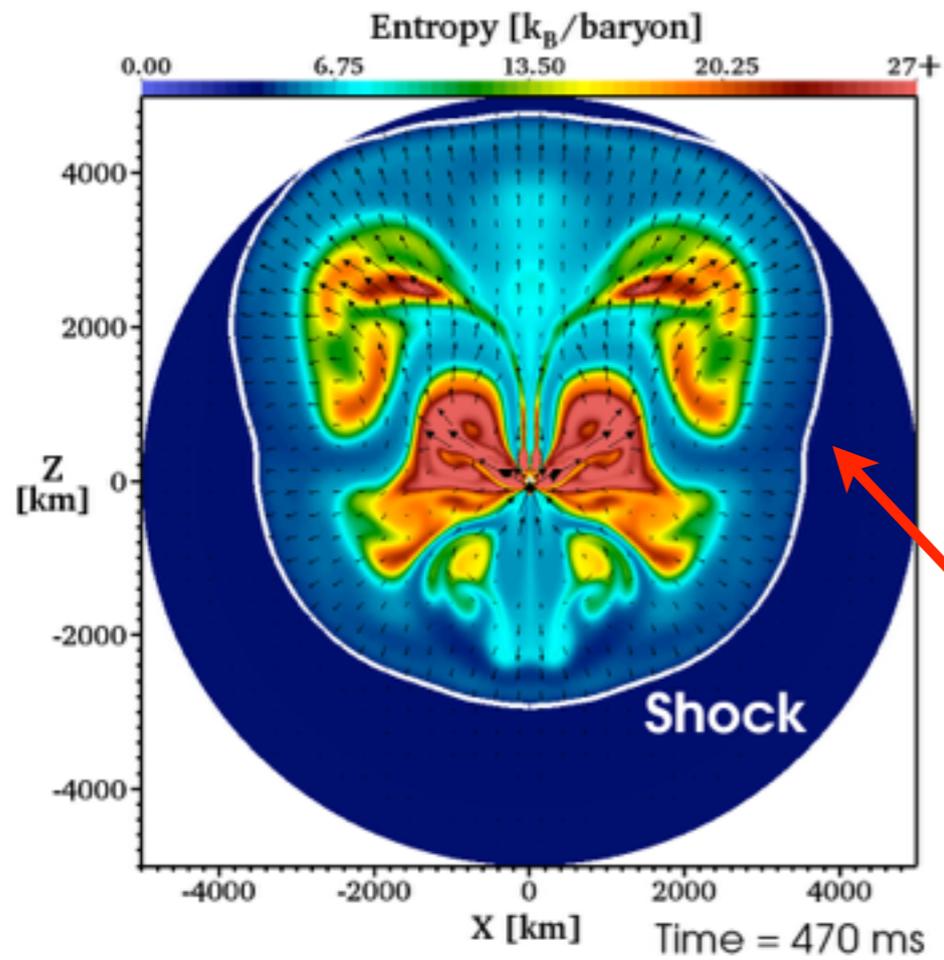
4. Pulsar Kicks.

Neutron Star Kicks



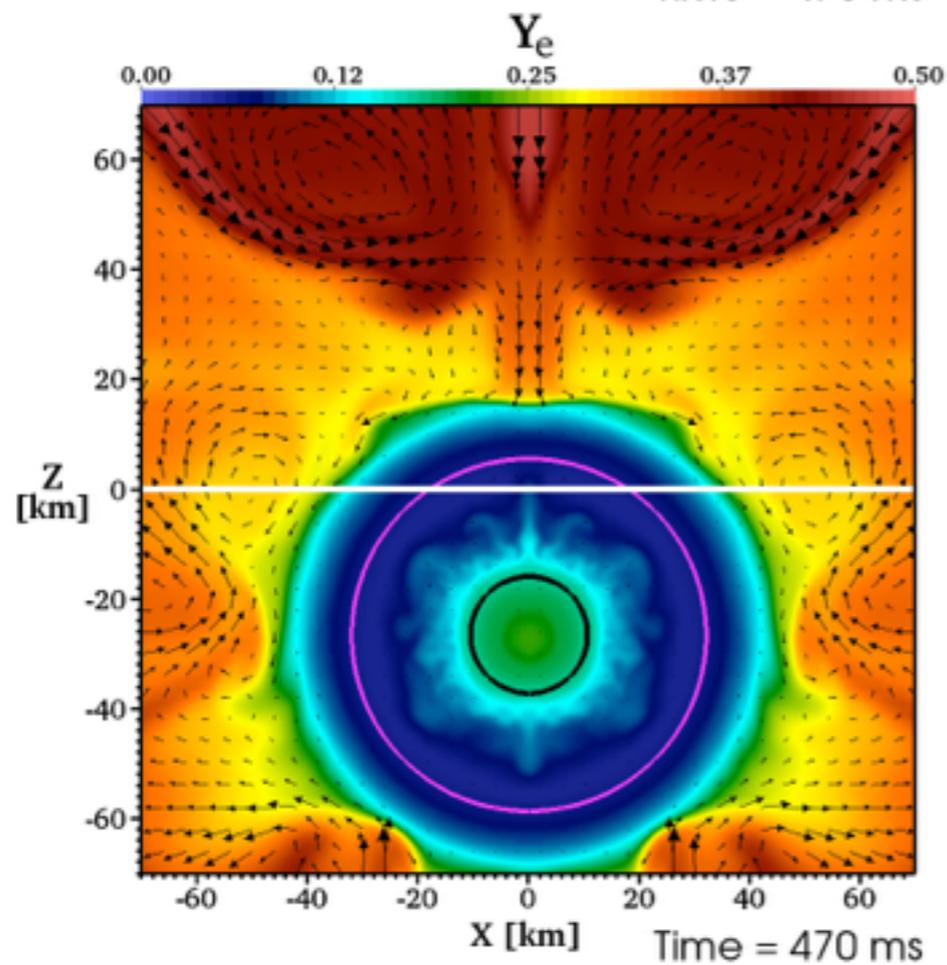
Nordhaus et al. 2010b

Neutron Star Kicks



Pulsar birth velocities typically 300 - 400 km s^{-1}

Explosion primarily in +Z direction...



...leads to NS recoil in -Z direction

Hydrodynamic Origin of Pulsar Kicks

Anisotropic neutrino emission
(neutrino “rockets”)
not important for kicks!

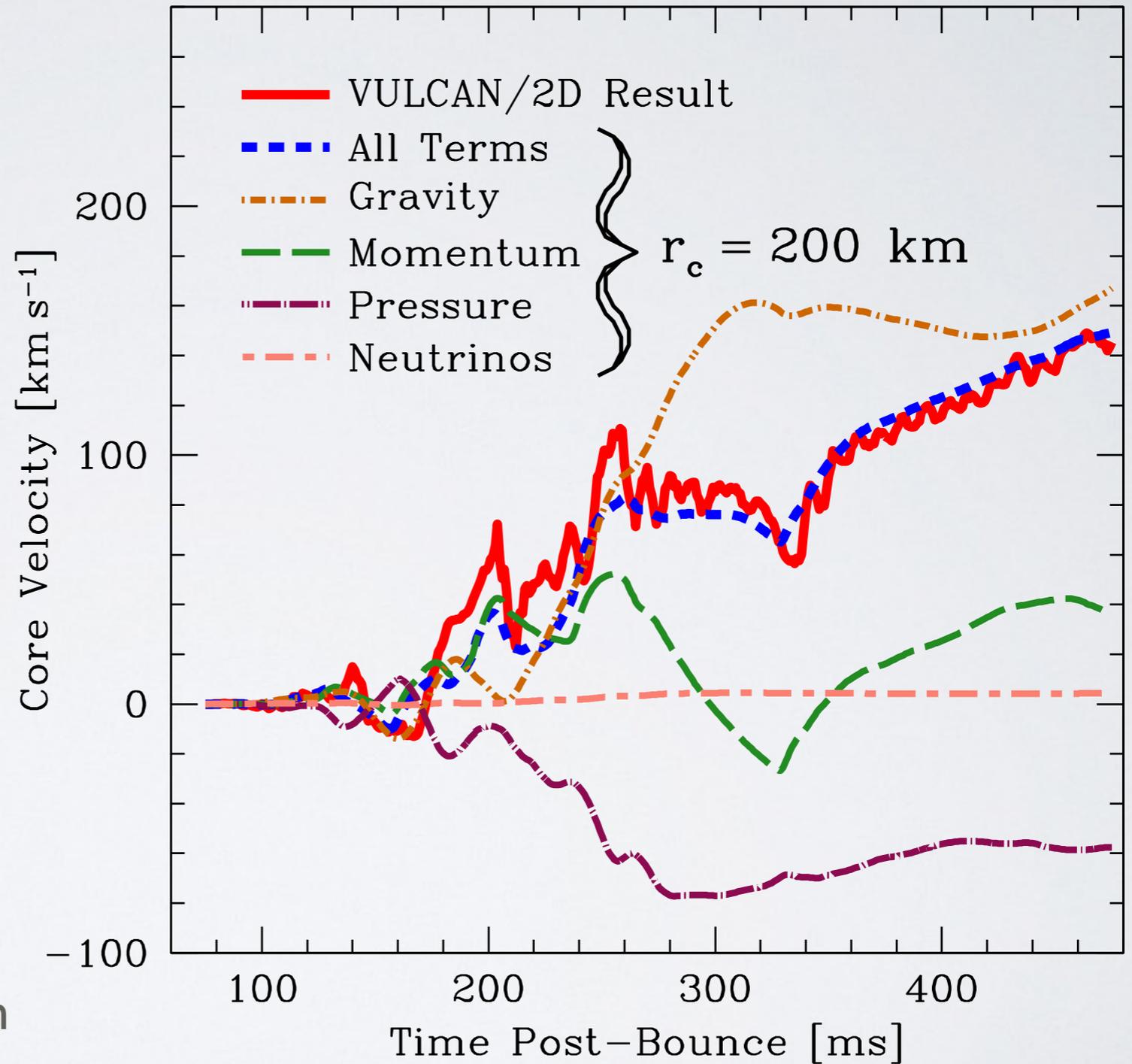
Gravitational tugboat effect
is important.

At end of simulation:

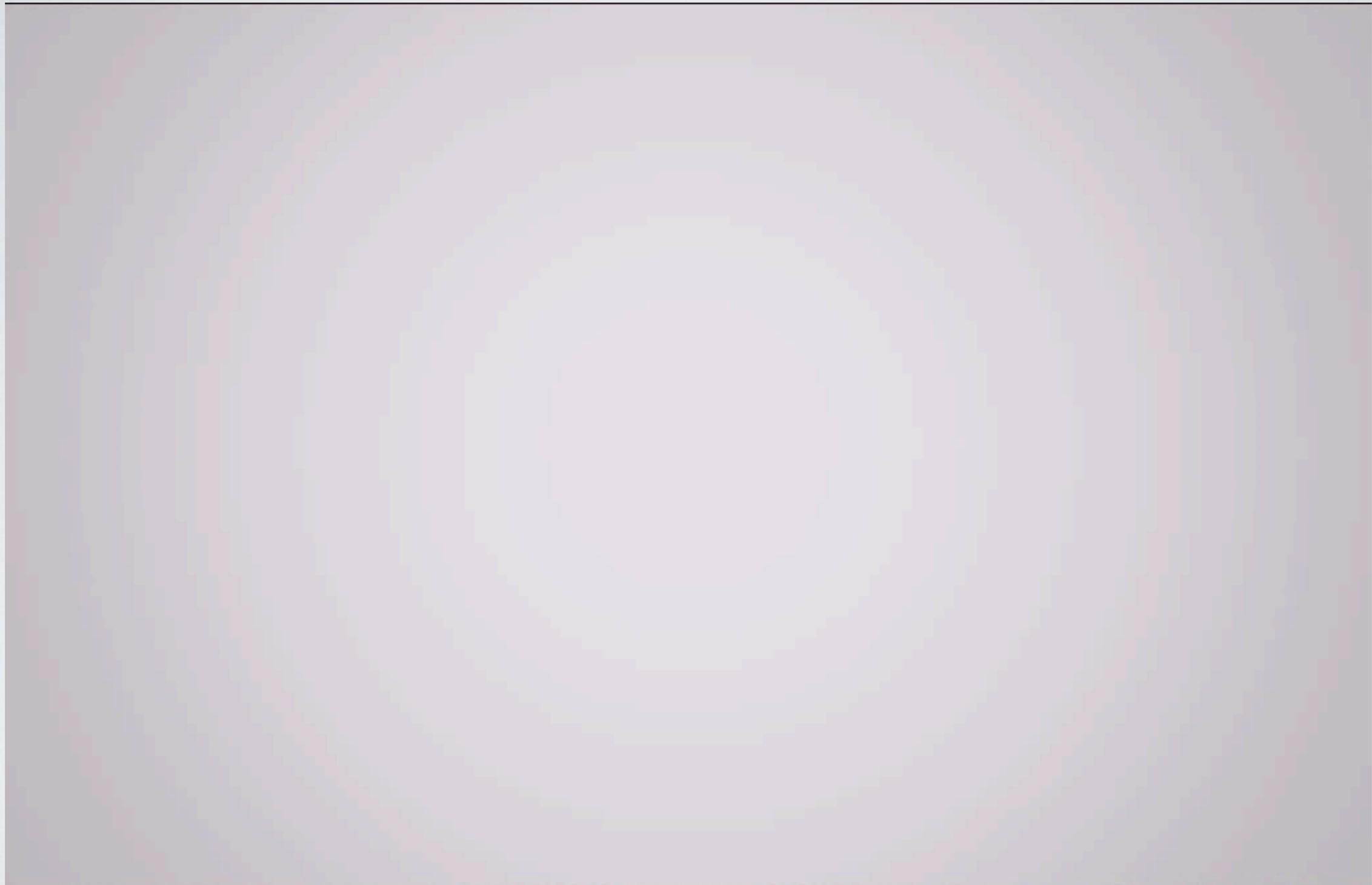
$$v_{\text{NS}} \sim 150 \text{ km s}^{-1}$$

$$a_{\text{NS}} \sim 350 \text{ km s}^{-2}$$

Requires $\sim 2\text{-}3$ seconds to reach
ballistic regime!



Radiation field is smooth...
... matter field is not.



Flux Vectors

ν_e
21.01 MeV

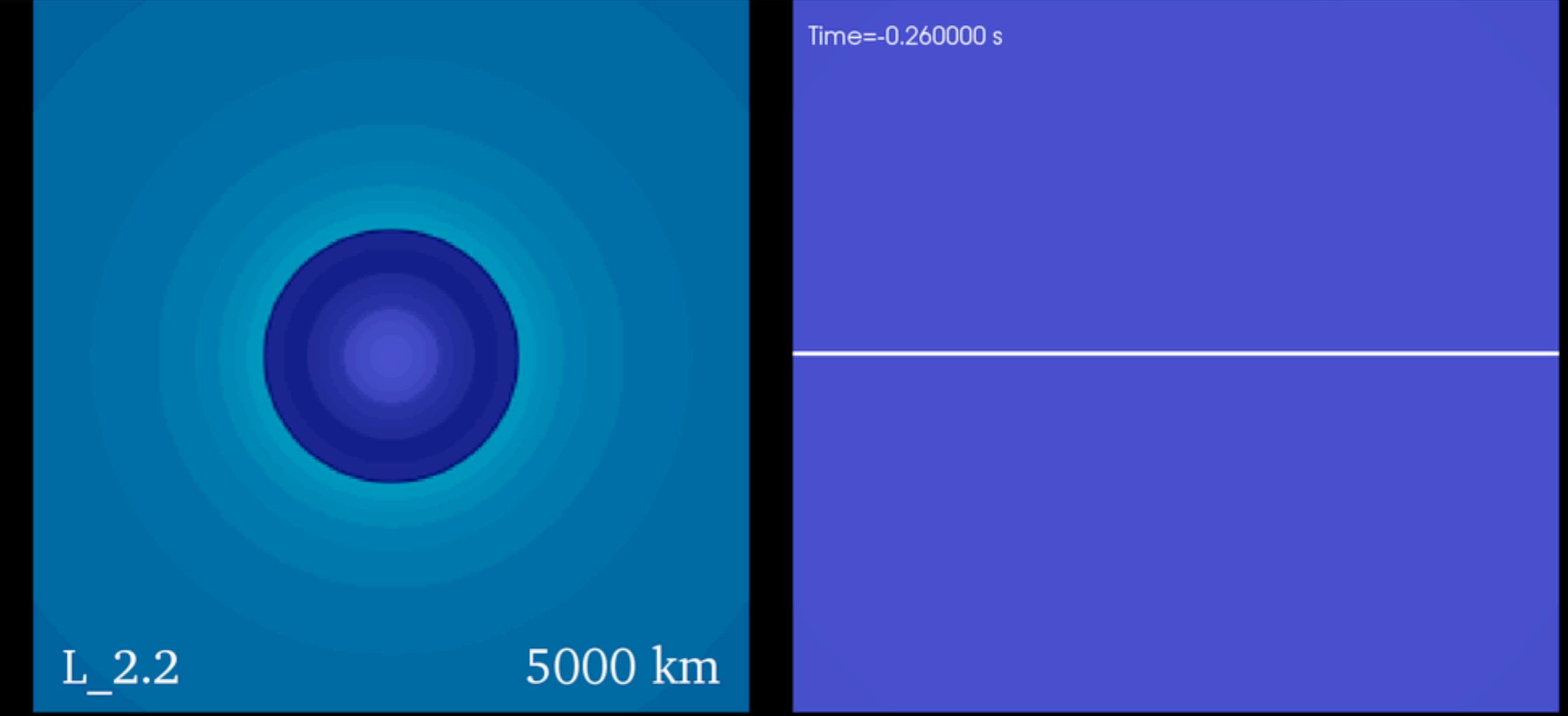
$\Omega = 0$

Time = -194.5 ms
Distance = 500.0 km

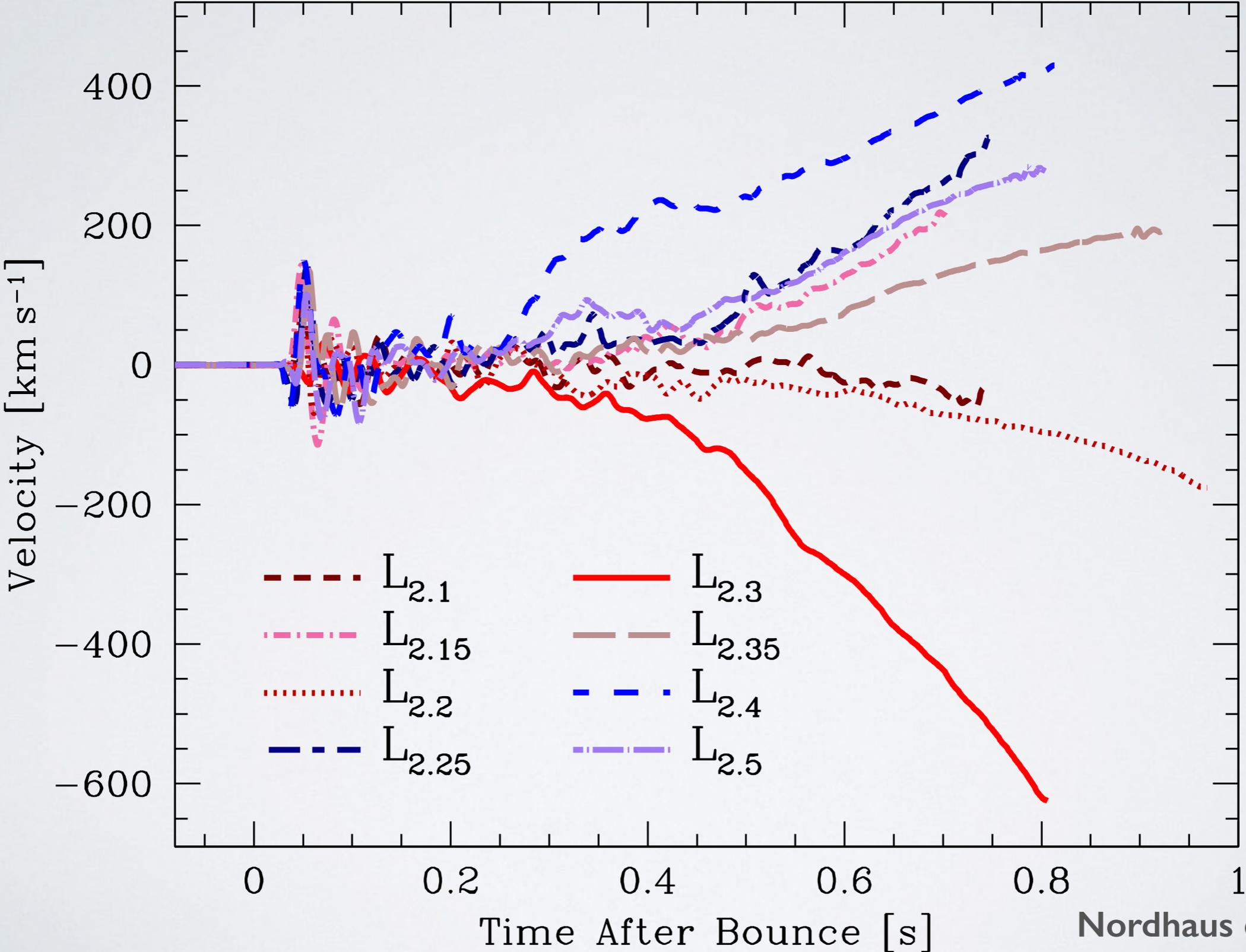


Hydrodynamic Origin of Pulsar Kicks

CASTRO - with neutrino heating/cooling scheme



Hydrodynamic Origin of Pulsar Kicks

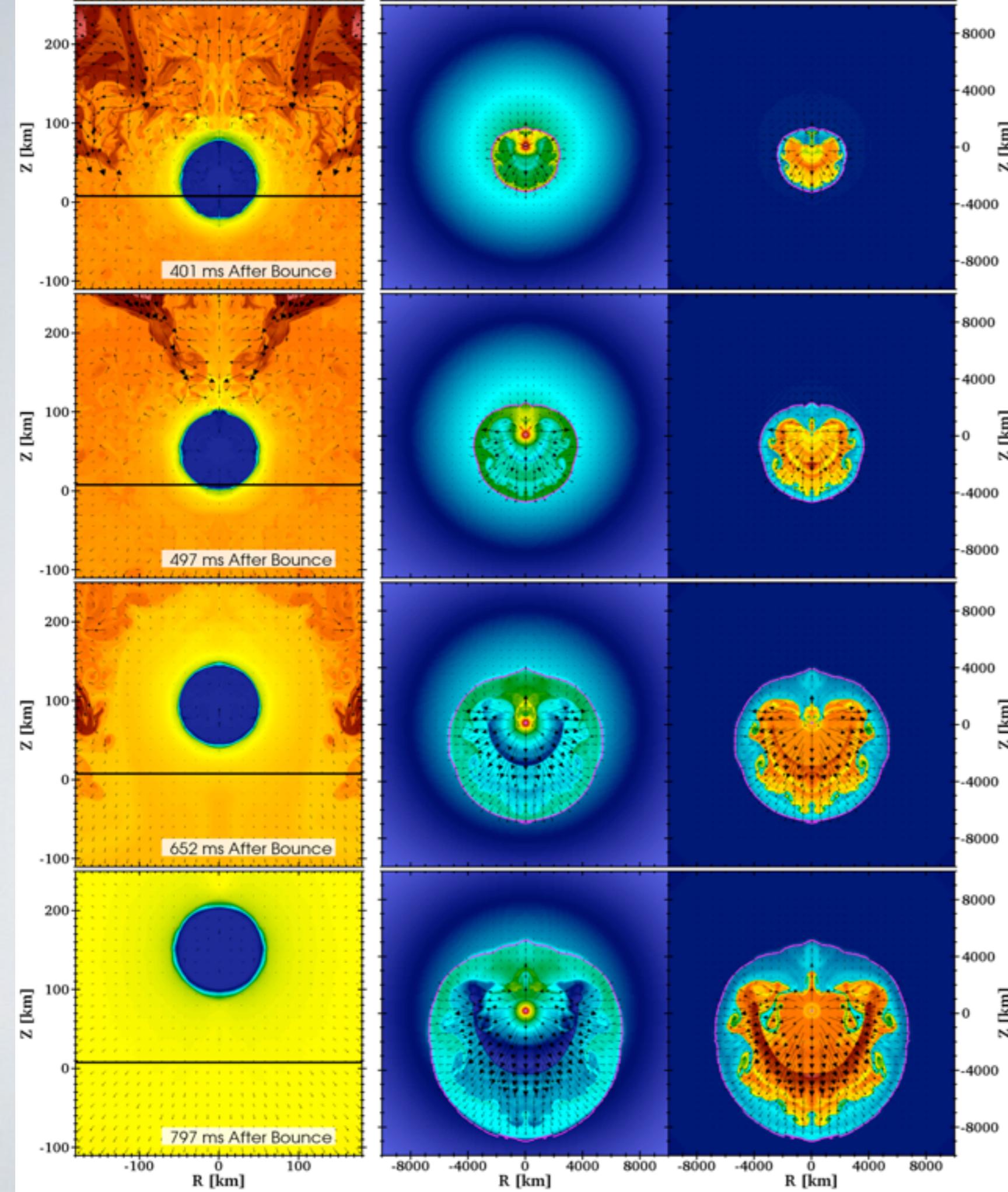


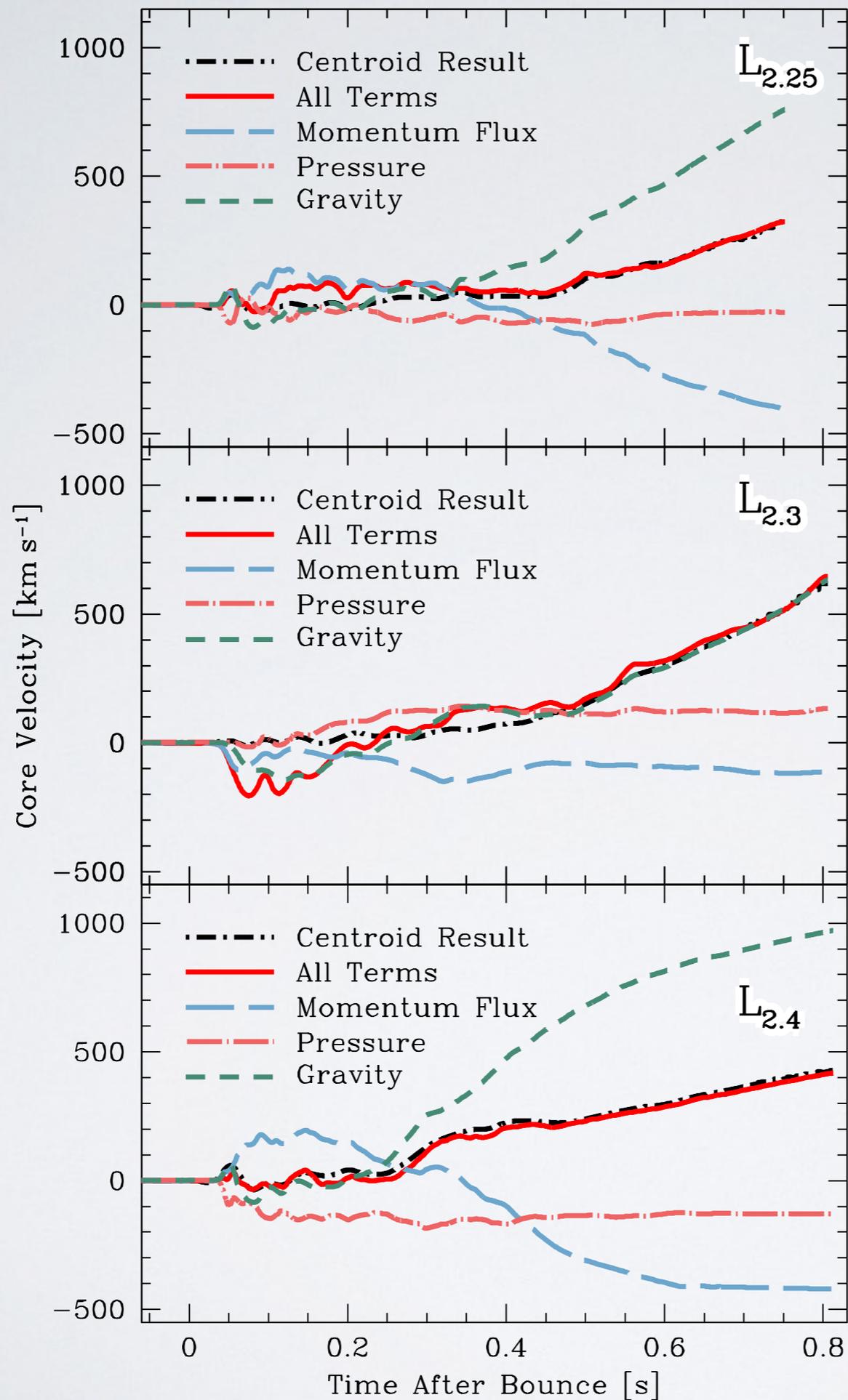
Pulsar Kicks

Gravitational effects are important.

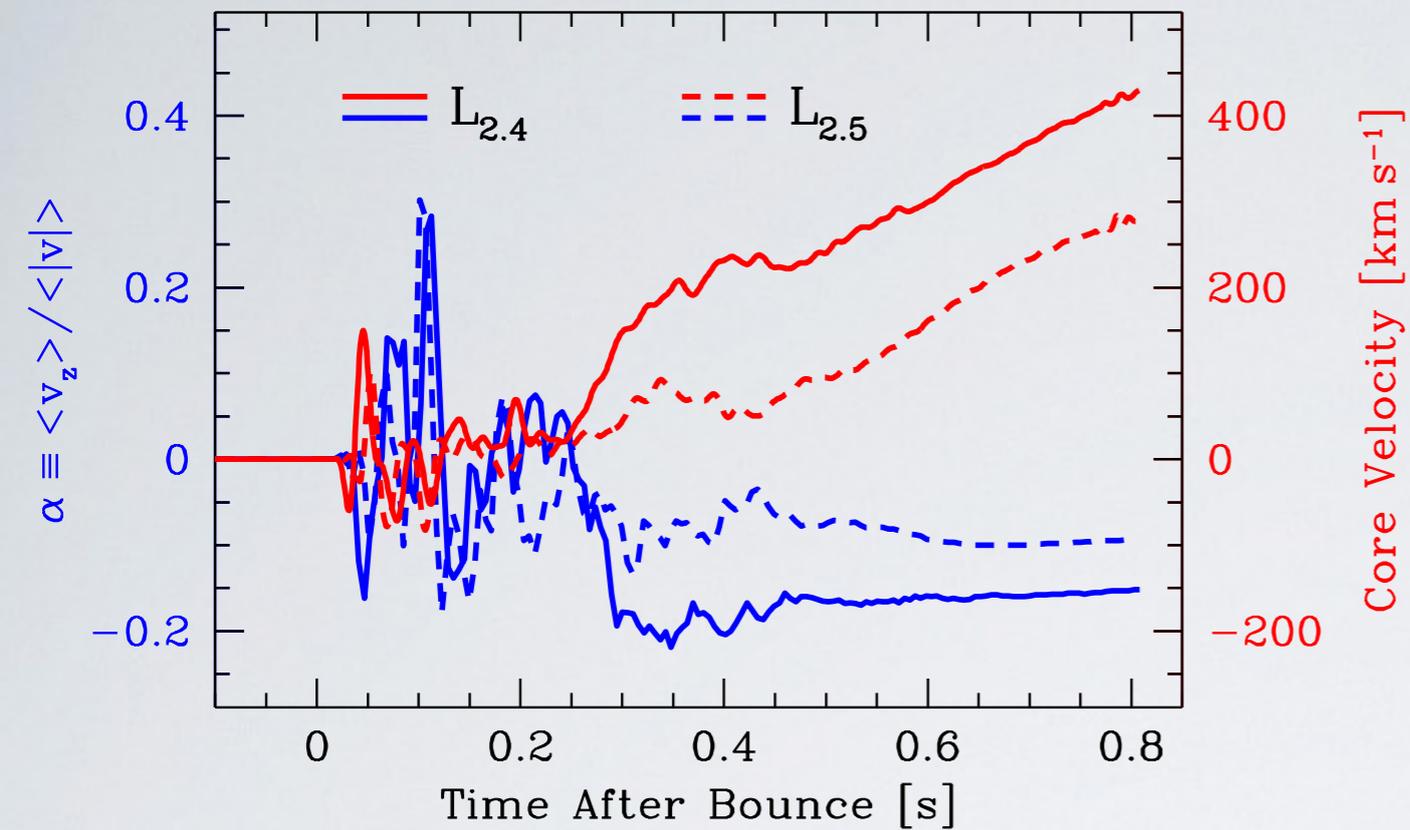
With AMR can follow evolution farther in time.

NS decoupled from surroundings





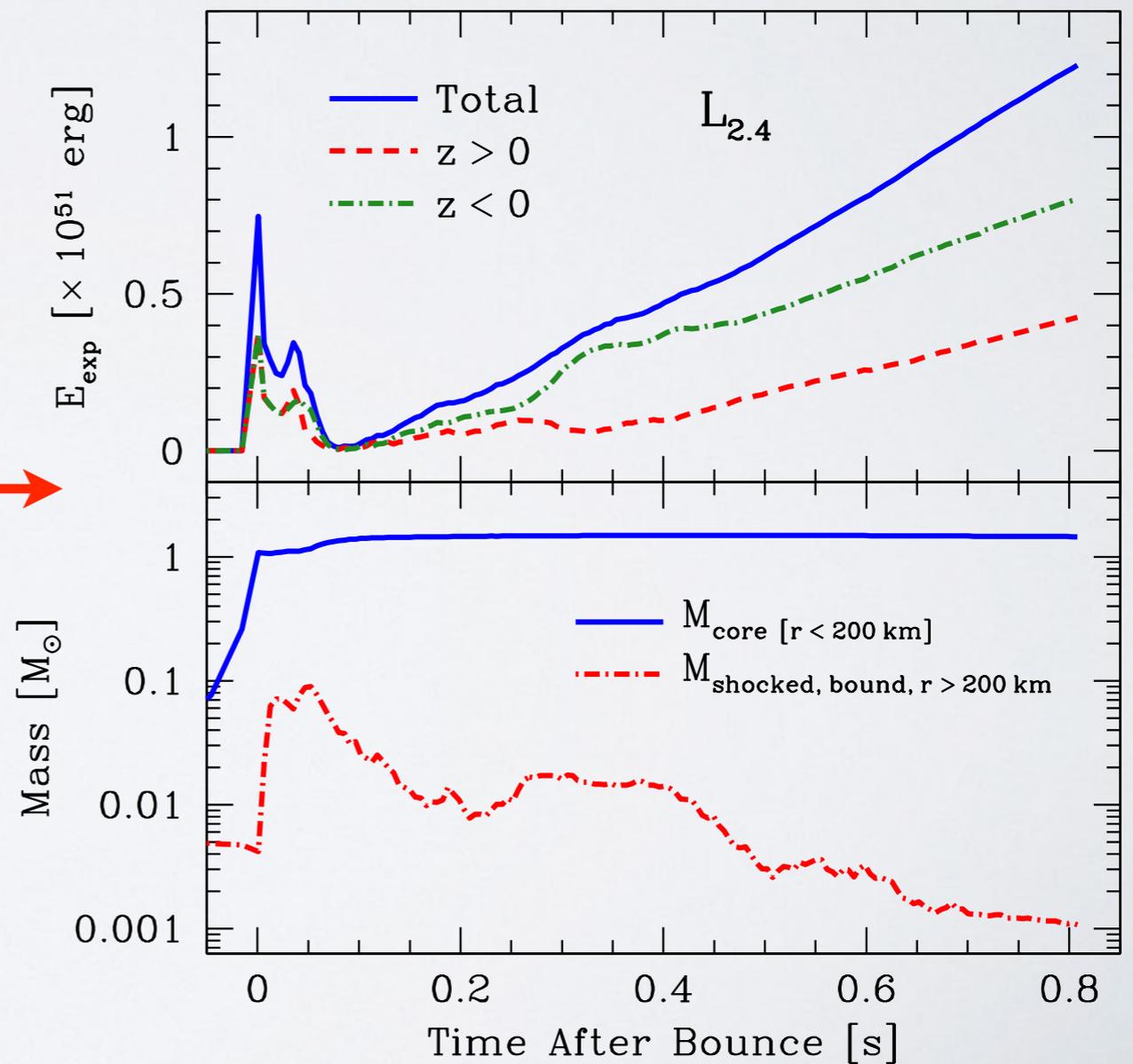
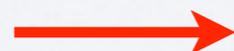
At late times, **gravity** of the **slow-moving ejecta** dominates.



The larger the degree of asymmetry, the larger the kick

Simulations achieve canonical supernova explosion energies.

Very little bound mass at end of the simulation.



Takeaways

State-of-the-art neutrino-driven core-collapse simulations do not reach supernova energies when they explode.

Inherent limitation of the neutrino-mechanism: Papish, Nordhaus, Soker 2015.

Alternate theories/processes needed: GR, MHD, nuclear burning.

Pulsar kicks are a hydrodynamic outcome of explosions.

