

Physics A300: Classical Mechanics I

Problem Set 7

Assigned 2002 October 28

Due 2002 November 4

Show your work on all problems!

1 Projectile Motion with Air Resistance (conclusion)

Recall that in problem 3 on Problem Set 6, you found that $y(t) = 0$ and that $x(t)$ and $z(t)$ obeyed the equations of motion

$$\ddot{x}(t) = -\frac{b}{m}\dot{x}(t) \quad (1.1a)$$

$$\ddot{z}(t) = -g - \frac{b}{m}\dot{z}(t) \quad (1.1b)$$

with initial conditions

$$x(0) = 0 \quad (1.2a)$$

$$z(0) = 0 \quad (1.2b)$$

and

$$\dot{x}(0) = v_0 \cos \alpha \quad (1.3a)$$

$$\dot{z}(0) = v_0 \sin \alpha \quad (1.3b)$$

- Solve the equations of motion (1.1) with initial conditions (1.3) to find $\dot{x}(t)$ and $\dot{z}(t)$ in terms of the time t , the coordinates $x(t)$ and $z(t)$, and the parameters of the problem.
- Solve the equations you got in part a) subject to the initial conditions (1.2) to find $x(t)$ and $z(t)$ and write the solution $\vec{r}(t)$ in terms of the time, the parameters of the problem, and the basis vectors \hat{x} , \hat{y} , \hat{z} .
- If T is the time when the projectile lands on level ground, write the equation which implicitly defines T in terms of the parameters of the problem.
- The result of part c) is called a *transcendental equation* because T appears both within and outside a transcendental function. It cannot be solved exactly, but when the air resistance is small, we can find an approximate solution perturbatively. As a first step in this process, define the dimensionless quantities $\beta = bv_0/mg$ and $\mathcal{T} = gT/v_0$ and use these definitions to rewrite your implicit equation for T as an equation relating only the dimensionless quantities \mathcal{T} , β , and α .

2 The Curl

- a) If $a(\vec{r})$ is a scalar field and $\vec{B}(\vec{r})$ is a vector field, show, by explicit evaluation of the left- and right-hand sides in Cartesian coordinates, that

$$\vec{\nabla} \times (a\vec{B}) = (\vec{\nabla}a) \times \vec{B} + a(\vec{\nabla} \times \vec{B}) . \quad (2.1)$$

- b) Writing the “del operator” in spherical coordinates according to Symon’s equation (3.124) allows us to write the curl of a vector as

$$\vec{\nabla} \times \vec{A} = \hat{r} \times \frac{\partial \vec{A}}{\partial r} + \frac{\hat{\theta}}{r} \times \frac{\partial \vec{A}}{\partial \theta} + \frac{\hat{\phi}}{r \sin \theta} \times \frac{\partial \vec{A}}{\partial \phi} . \quad (2.2)$$

Use this, along with Symon’s equation (3.99), to calculate i) $\vec{\nabla} \times \hat{r}$; ii) $\vec{\nabla} \times \hat{\theta}$; iii) $\vec{\nabla} \times \hat{\phi}$.

- c) Using the results of parts a) and b), and writing a vector field $\vec{A}(\vec{r})$ as

$$\vec{A}(\vec{r}) = A_r(r, \theta, \phi) \hat{r} + A_\theta(r, \theta, \phi) \hat{\theta} + A_\phi(r, \theta, \phi) \hat{\phi} \quad (2.3)$$

show that the curl in spherical coordinates is

$$\begin{aligned} \vec{\nabla} \times \vec{A} = & \left(\frac{1}{r} \partial_\theta A_\phi - \frac{1}{r \sin \theta} \partial_\phi A_\theta + \frac{\cos \theta}{r \sin \theta} A_\phi \right) \hat{r} + \left(\frac{1}{r \sin \theta} \partial_\phi A_r - \partial_r A_\phi - \frac{1}{r} A_\phi \right) \hat{\theta} \\ & + \left(\partial_r A_\theta - \frac{1}{r} \partial_\theta A_r + \frac{1}{r} A_\theta \right) \hat{\phi} \end{aligned} \quad (2.4)$$

3 Force, Potential and Torque

Consider the force field

$$\vec{F}(\vec{r}) = V_0 \frac{x \hat{x} + y \hat{y}}{x^2 + y^2} \quad (3.1)$$

- a) By explicitly calculating the (three-dimensional) curl $\vec{\nabla} \times \vec{F}$, verify that this is a conservative force.
- b) Invert Symon’s equation (3.89) to obtain expressions for \hat{x} , \hat{y} and \hat{z} in terms of the cylindrical coordinates ρ , ϕ and z and the basis vectors $\hat{\rho}$, $\hat{\phi}$, and \hat{z} . Simplify your answer as much as possible.
- c) Use Symon’s equation (3.87) and the results of part b) to write \vec{F} above entirely in terms of the cylindrical coordinates ρ , ϕ and z and the basis vectors $\hat{\rho}$, $\hat{\phi}$, and \hat{z} (and the constant V_0). Simplify your answer as much as possible.
- d) Working in cylindrical coordinates, find the potential energy $V(\rho, \phi, z)$ such that $\vec{F} = -\vec{\nabla}V$. Include in your result an arbitrary constant (so that you capture the entire family of possible potentials) and indicate its units.
- e) Calculate the vector torque \vec{N} due to this force (in either Cartesian or cylindrical coordinates), and verify that the torque about the z axis vanishes.